Plasma and Dry Micro/Nanotechnologies 3. Chemical and Plasma Kinetics

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• 3. Chemical and Plasma Kinetics

- 3.1 Chemical Kinetics
- 3.2 Plasma Kinetics
- 3.3 Elementary Processes in Plasma Involving Atoms
- 3.4 Elementary Processes in Plasma Involving Molecules

Chemical vapor deposition (CVD): based on chemical reactions. Plasmachemical processing of materials (plasma treatment in reactive gases, plasma etching, plasma enhanced CVD - PECVD): chemical reactions with reactive species activated by hot electrons.

- homogeneous if reactants are of the same phase (e.g. reactions in gas phase)
- heterogenous if reactants are of two or more different phases (e.g. reaction between gas and solid - deposition, etching)

Related terms:

- reaction rate, reaction order, rate constant, Arrhenius plot
- ellastic collisions (Coulomb, polarization scattering with permanent or induced dipole, hard sphere scattering)
- ▶ inellastic collisions (ionization, recombination, excitation, dissociation ...)
- collisional cross-section and its relation to the reaction constant

3.1 Chemical Kinetics

Consider the reaction

$$aA + bB \cdots \longrightarrow qQ + sS \dots$$
 (1)

The reactants A and B react with the rate

$$r'_{A} = -\frac{\mathrm{d}N_{A}}{\mathrm{d}t}$$
 $r'_{B} = -\frac{\mathrm{d}N_{B}}{\mathrm{d}t}.$ (2)

The products Q and S are created with the rate

$$r'_{S} = \frac{\mathrm{d}N_{S}}{\mathrm{d}t} \quad r'_{Q} = \frac{\mathrm{d}N_{Q}}{\mathrm{d}t},\tag{3}$$

where N_i is molar amount (usually in kmol) and t is time.

$$r' = -\frac{1}{a}\frac{\mathrm{d}N_A}{\mathrm{d}t} = -\frac{1}{b}\frac{\mathrm{d}N_B}{\mathrm{d}t} = \frac{1}{q}\frac{\mathrm{d}N_Q}{\mathrm{d}t} = \frac{1}{s}\frac{\mathrm{d}N_S}{\mathrm{d}t}.$$
(4)

The reaction rate is generally expressed on an intensive basis, e.g. reaction volume V.

$$r = \frac{-1}{aV} \frac{dN_A}{dt} = \frac{-1}{aV} \frac{d}{dt} (C_A V) = \frac{-1}{aV} \left(V \frac{dC_A}{dt} + C_A \frac{dV}{dt} \right),$$
(5)

where C_A is the molar concentration of A (kmol/m³). When the volume is constant

$$r = \frac{-1}{a} \frac{dC_A}{dt}.$$
 (6)

The reaction rate can be expressed also as

$$r = k_{AB} C^a_A C^b_B. \tag{7}$$

where k_{AB} is the rate constant of the reaction $aA + bB \rightarrow$ products.

- The powers a, b are called partial orders of the reactions.
- The sum a + b is (global) order.

The partial orders should be in fact replaced by a', b' determined experimentally because they coincide with the stoichiometric coefficients only for elementary processes.

Arrhenius concluded from thermodynamics (T is temperature, R and $k_{\rm B}$ gas and Boltzmann constants):

$$\ln k = -\frac{\varepsilon_{\rm a}}{RT} + \ln A_0 = -\frac{E_{\rm a}}{k_{\rm B}T} + \ln A_0$$

 $\varepsilon_{\rm a}$ and $E_{\rm a}$ activation energies in J/mol and J, A_0 constant frequency factor.

3.1 Chemical Kinetics

For non-reversible firts order reaction $A \longrightarrow Q$, we have at constant volume

r

$$r = -\frac{dC_A}{dt} = kC_A.$$
 (8)

If we know the rate constant k (hr⁻¹), we can calculated the reaction rate r for any concentration of the reactant.

The rate constant *k* can be calculated from Eq. (8) if we know the time change of the concentration \Rightarrow differential method. Integral method uses integrated Eq. (8):

$$\ln\left(\frac{C_{A0}}{C_A}\right) = kt. \tag{9}$$

The semilog plot of C_{A0}/C_A in dependence on time gives *k*.

Now, think about the expressions for

- reversible 1st order reaction $A \simeq_2^1 Q$
- irreversible 2nd order reaction $2A \rightarrow Q + S$
- irreversible 2nd order reaction $A + B \rightarrow Q + S$

We can encounter also with **complex reactions**: parallel, consequtive, free radical chain reaction (initiation, propagation, termination), free radical addition

3.2 Plasma Kinetics

Elastic collisions (momentum and energy are conserved):

- Coulomb collisions between two charged particles (e-e, e-ion, ion-ion)
- > polarization scattering with **permanent dipole** (for molecules with permanent dipole)
- polarization scattering with induced dipole (e-neutral for electrons with low energies, ion-neutral collisions)
- ▶ hard sphere between neutrals, e-neutral for very low electron energies (approx.)

Process	U(r)	σ	ν or K
Coulomb	1/r	$1/v_R^4$	$1/v_{R}^{3}$
Permanent dipole	$1/r^2$	$1/v_R^2$	$1/v_{\rm R}$
Induced dipole	$1/r^4$	$1/v_R$	const
Hard sphere	$1/r^i, i \to \infty$	const	$v_{\mathbb{R}}$

TABLE 3.1. Scaling of Cross Section σ , Interaction Frequency ν , and Rate Constant K, With Relative Velocity ν_{R} , for Various Scattering Potentials U

after Lieberman & Lichtenberg 1994

Inelastic collisions: ionization, recombination, excitation, dissociation

3.2 Plasma Kinetics

Electrons in plasma gain high energies (in the order of 1-10 eV) due to acceleration by electric field.

Since electrons collide with heavy particles (atoms, molecules) they change direction of their velocity or even loose the energy.

Collisions between electrons and heavy particles (according to the electron energy $E_{\rm e}$):

- E_e < 2 eV (depending on the atom/molecule): elastic collisions with very small fractional energy transfer (see slide in Atomic Collisions).
- ▶ 2 eV $< E_{\rm e} < 15$ eV (approx.): variety of inelastic collisions $\Rightarrow E_{\rm e}$ is partially converted into internal energy of the target molecule (atom)
- $E_{\rm e} > 15 \, {\rm eV}$ (approx.): ionization (sustains the discharge)

Rate constant *k* for reaction of two particles with velocities \vec{v}_1 , \vec{v}_2 can be calculated from cross section σ

$$k = \langle \sigma(\mathbf{v}_{\mathrm{R}})\mathbf{v}_{\mathrm{R}}\rangle_{\mathbf{v}_{1},\mathbf{v}_{2}} = \int \sigma(\mathbf{v}_{\mathrm{R}})\mathbf{v}_{\mathrm{R}}f_{1}(\vec{\mathbf{v}}_{1})f_{2}(\vec{\mathbf{v}}_{2})d\mathbf{v}_{1}^{3}d\mathbf{v}_{2}^{3}$$

where $v_{\rm R} = |\vec{v}_1 - \vec{v}_2|$ and $f_1(\vec{v}_1)$, $f_2(\vec{v}_2)$ are velocity distribution functions.

Plasma Kinetics

The velocity distributions are taken isotropic Maxwellian.

$$f(v) = \left(\frac{m}{2\pi k_{\rm B}T}\right)^{3/2} \exp\left(-\frac{mv^2}{2k_{\rm B}T}\right)$$

where m and T are particle mass and temperature.

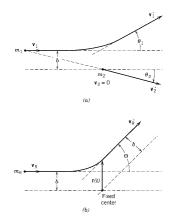
If the characteristics velocities of target particles are much less than those of incident particles (e.g. electron collision with heavy particle) $v_{\rm R} \approx |\vec{v}_1| \equiv v$.

$$k(T) = \langle \sigma(v)v \rangle_{v} = \left(\frac{m}{2\pi k_{\rm B}T}\right)^{3/2} \int_{0}^{\infty} \sigma(v)v \exp\left(-\frac{mv^{2}}{2k_{\rm B}T}\right) 4\pi v^{2} dv$$

If we consider collision of two different heavy particles

$$k(T) = \langle \sigma(\mathbf{v}_{\mathrm{R}})\mathbf{v}_{\mathrm{R}} \rangle_{\mathbf{v}_{1},\mathbf{v}_{2}} = \left(\frac{m_{\mathrm{R}}}{2\pi k_{\mathrm{B}}T}\right)^{3/2} \int_{0}^{\infty} \sigma(\mathbf{v}_{\mathrm{R}})\mathbf{v}_{\mathrm{R}} \exp\left(-\frac{mv_{\mathrm{R}}^{2}}{2k_{\mathrm{B}}T}\right) 4\pi v_{\mathrm{R}}^{2} \mathsf{d}v_{\mathrm{R}}$$

where $m_{\rm R}$ is reduce mass.



Scattering in (a) laboratory system, (b) the center of mass (CM) system (after Lieberman & Lichtenberg 1994).

for $m_1=m_{
m e},\,m_2=M$ and $m_{
m e}\ll M$

Electron - atom elastic collision:

- momentum and energy are conserved,
- treated as hard-sphere scattering

Fraction of energy lost by the projectile in the laboratory system

$$\gamma = rac{E_{
m t}}{E_{
m i}} = rac{4m_{
m 1}m_{
m 2}}{(m_{
m 1}+m_{
m 2})^2}\cos^2 heta_{
m 2}$$

and in the CM system

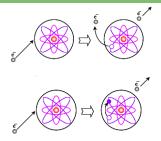
$$\gamma = rac{E_{
m t}}{E_{
m i}} = rac{2m_{
m 1}m_{
m 2}}{(m_{
m 1}+m_{
m 2})^2}(1-\cos\Theta)$$

Average loss obtained by averaging over all angles Θ using differential cross section $\sigma(\nu_{\rm R},\Theta)$ as distribution function

$$\langle \gamma \rangle_{\Theta} = \frac{2m_1m_2}{(m_1 + m_2)^2} \frac{\int_0^{\pi} (1 - \cos\Theta)\sigma(v_{\rm R}) 2\pi \sin\Theta {\rm d}\Theta}{\int_0^{\pi} \sigma(v_{\rm R}) 2\pi \sin\Theta {\rm d}\Theta}$$

$$\langle M \rangle_{\Theta} = \frac{2m_1m_2}{(m_1+m_2)^2} \approx \frac{2m_{\rm e}}{M} \approx 10^{-4}$$

3.3 Atomic Processes / Excitation Processes



Electron-metastable ionization

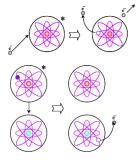
Electron impact ionization

$$\mathrm{e^-} + \mathrm{A} \longrightarrow \mathrm{e^-} + \mathrm{e^-} + \mathrm{A^+}$$

Electron impact excitation

$$e^- + A \longrightarrow e^- + A^*$$

 $\rm A^*$ can have quite different chemical reactivity towards the surface. Some excited atoms have very long lifetimes (\approx 1–10 ms) \Rightarrow metastables



 $e^- + A^* \longrightarrow e^- + e^- + A^+$

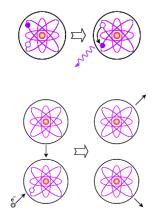
Since the metastable atom is already excited, less energy is required.

Metastable-neutral ionization

$$\mathrm{A}^* + \mathrm{B} \longrightarrow \mathrm{A} + \mathrm{e}^- + \mathrm{B}^+$$

If the ionization energy of the neutral B is less than the excitation energy of the metastable $A^* \Rightarrow$ **Penning** ionization (He^{*} 19.8, Ne^{*} 16.7, Ar^{*} 11.7 eV)

Atomic Processes - Relaxation and Recombination



De-excitation

 $\mathbf{A}^* \longrightarrow \mathbf{A} + \mathbf{h} \nu$

In most cases, the relaxation of electronically excited states is practically instantaneous (\approx 10 ns).

Electron-ion recombination

 $e^- + A^+(+C) \longrightarrow A^*(+C)$

A third-body (neutrals, reactor walls) must be involved to conserve energy and momentum.

Radiative recombination

$$e^- + A^+ (+C) \longrightarrow A + h\nu (+C)$$

Electron attachment

$$e^- + A (+C) \longrightarrow A^- (+C)$$

Ion-ion recombination

$$A^+ + A^- \longrightarrow A + A$$

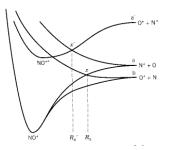
Heavy-Particle Collisions - Charge Transfer

- ► In general, the energy level from which e^- is released is not equal to the energy level into which the electron is captured \Rightarrow energy defect ΔW .
- For $\Delta W \neq 0$, the kinetic energy of the colliding particles is not conserved in the collision.

Resonant charge transfer If atom and ion are parent and child, the transfer occur with $\Delta W = 0$

$$A^+(fast) + A(slow) \longrightarrow A(fast) + A^+(slow)$$

Cross section is larger for low energies, important process in weakly ionized plasmas.



Nonresonant charge transfer

$$A^+ + B \longrightarrow A + B^+$$

lllustrated for $\rm N^+ + O$ and $\rm O^+ + N$ (ioniz. potential of N and O are 14.53 and 13.61 eV, respectively) Exothermic reaction a-x-b

 $\rm N^+ + O \longrightarrow N + O^+$

does not have a threshold energy, products share an increased kinetic energy of 0.92 eV.

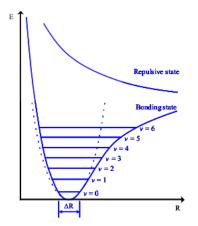
The inverse endothermic reaction ($\varepsilon_{\rm thr}=0.92\,\text{eV})$ has very small rate constant at thermal energies

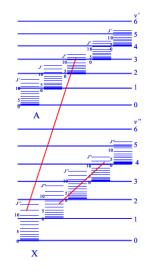
 $O^+ + N \longrightarrow O + N^+$

but if ${\rm O^+}$ or N are excited, the reaction a'-x'-a has no $\varepsilon_{\rm thr}$ and σ can be large at thermal energies.

3.4 Elementary Processes in Plasma Involving Molecules

In molecules, excitation of vibrational and rotational states (besides electronic states) are possible:





Electron Collisions with Molecules - Interaction Times

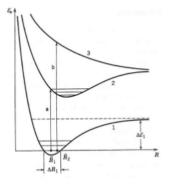
Interaction times:

- electron motion in the molecule $t_{\rm at} \approx 10^{-16}$ – 10^{-15} s
- ▶ interaction time of e⁻ with molecule $t_{int} \approx 10^{-16}$ -10⁻¹⁵ s
- ► typical vibrational period of molecule $t_{\rm vib} \approx 10^{-14}$ -10⁻¹³ s
- ▶ typical time for molecule dissociation $t_{\rm diss} \approx 10^{-14} 10^{-13} \, {\rm s}$
- ► typical transition time for electric dipole radiation $t_{\rm rad} \approx 10^{-9}$ -10⁻⁸ s
- ▶ typical time between e-molecule collisions in a low pressure plasma $t_{\rm col} \approx 10^{-8}$ –10⁻⁶ s

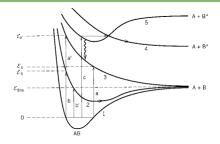
 $t_{\rm at} \sim t_{\rm int} \ll t_{\rm vib} \sim t_{\rm diss} \ll t_{\rm rad} < t_{\rm col}$

For electron collisional excitation of molecule to an excited electronic state, the new vibrational (and rotational) state can be determined by freezing the nuclear motion during the collision $t_{\rm int} \ll t_{\rm vib} \sim t_{\rm diss} \Rightarrow$ Frank-Condon principle

Since $t_{\rm diss} \ll t_{\rm rad},$ if energetics permit, the molecule will dissociate instead of de-exciting.



Electron Collisions with Molecules - Dissociation



Dissociation cross section rises linearly from threshold $\varepsilon_{\rm thr} \approx \varepsilon_1$ to a max. value (typically 10⁻¹⁵ cm²) at ε_2 and then falls off as $1/\varepsilon$:

$$\begin{aligned} \sigma_{\rm diss} &= 0 \quad \varepsilon < \varepsilon_1 \\ \sigma_{\rm diss} &= \sigma_0 \frac{\varepsilon - \varepsilon_1}{\varepsilon_1} \quad \varepsilon_1 < \varepsilon < \varepsilon_2 \\ \sigma_{\rm diss} &= \sigma_0 \frac{\varepsilon_2 - \varepsilon_1}{\varepsilon} \quad \varepsilon > \varepsilon_2 \\ \sigma_0 &= \pi \left(\frac{e}{4\pi\varepsilon_0\varepsilon_1}\right)^2 \end{aligned}$$

Dissociation

key role for plasma chemistry of low pressure discharges:

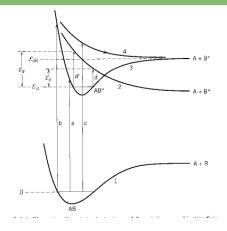
 $e^- + AB \longrightarrow A + B + e^-$

Collisions *a* and *a'*: ground state v = 0excited to repulsive state of AB, energy $\langle \varepsilon_a - \varepsilon_{diss}, \varepsilon_{a'} - \varepsilon_{diss} \rangle$ shared among the dissociation products A and B. Typically, $\varepsilon_a - \varepsilon_{diss} \approx$ few eV \Rightarrow hot neutral fragments (profound effect on plasma chemistry of growing films if hitting the substrate surface)

Collisions *b* and *b'*: ground state excited to an attractive state of AB but energy exceeds $\varepsilon_{\rm diss} \Rightarrow$ dissociation of AB resulting in fragments having energies from thermal up to $\varepsilon_{\rm b} - \varepsilon_{\rm diss} \approx$ few eV.

Collision c: excitation to bound state AB* that radiates creating A + B or $AB^*(bound) \longrightarrow AB^*(unbound) \longrightarrow A+B^*$

Electron Collisions with Molecules - contin.



Dissociative Ionization (in addition to normal ionization)

 $e^- + AB \longrightarrow A + B^+ + e^-$

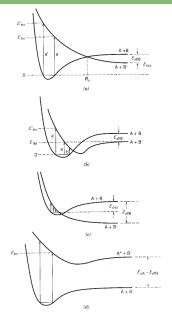
is common for polyatomic molecules. Formation of molecular ion (collision *a*) for threshold energy ε_{iz} . Collisions *b*, *c* for higher threshold energies $\varepsilon_{diz} \Rightarrow$ fast ion and neutral.

Dissociative Recombination

 $\mathrm{e^-} + \mathrm{AB^+} \longrightarrow \mathrm{A} + \mathrm{B^*}$

collisions d, $d' \Rightarrow$ fast excited neutral fragments.

Electron Collisions with Molecules - contin.



Dissociative Electron Attachment

 $e^- + AB \longrightarrow A + B^-$

important in discharges containing atoms with positive electron affinities because of

- production of negative ions
- threshold energy for dissociation is generally lower than for pure dissociation processes

(a) e^- capture into repulsive state \Rightarrow autodetachment or dissociation; autodetach. rate $\sqrt{M_{\rm R}/m} \approx 100 \times$ dissoc. rate ($M_{\rm R}$ reduced mass); hot fragments (b) AB⁻ bound state \Rightarrow *a*, *a'* dissociative attach. with low energy fragments; *b* collision AB^{-*} \rightarrow e⁻ + AB (c) for few molecules (e.g. halogens) $\varepsilon_{\rm affB} > \varepsilon_{\rm diss} \Rightarrow$ slow e⁻ produce hot A + B⁻; max. σ as high as 10⁻¹⁶ cm²

Polar Dissociation (d)

 $e^- + AB \longrightarrow A^+ + B^- + e^-$

- Maximum cross section and its dependence on electron impact energy are similar to pure dissociation.
- Threshold energy is generally large.

Electron Collisions with Molecules - contin.

Electron Impact Detachment

 $e^- + AB^- \longrightarrow AB + 2e^-$

similar to electron-neutral ionization with el. affinity $\varepsilon_{\rm aff}$ of AB playing the role of the ionization potential BUT the peak in cross section is shifted to energies of 10–20 $\varepsilon_{\rm aff}$ due to repulsive Coulomb force between e^- and AB^- .

Vibrational and Rotational Excitations

Typically it is a two step process:

$$e^- + AB(v = 0) \longrightarrow AB^-$$

Life time of AB^- is $10^{-15}-10^{-10}$ s,

i. e. comparable or larger than its vibrational timescale $10^{-14} \, s \Rightarrow$

$$AB^- \longrightarrow AB(v > 0) + e^-$$

Complex reaction schemes for O_2 plasma - 2^{nd} order reactions

Number	Reaction	Rate Constant (cm ³ /s)
Reactions an	tong e, O2, O2, O, and O	
1	e + O2 momentum transfer	4.7E-8T _e ^{0.5}
2	$e + O_2 \rightarrow O^- + O$	$1.07E - 9T_e^{-1.391} \exp(-6.26/T_e)$
3	$e + O_2 \rightarrow 2O + e$	6.86E-9exp(-6.29/T _e)
4	$e + O_2 \rightarrow O_2^+ + 2e$	2.34E-9Te ^{1.03} exp(-12.29/Te)
5	$e + O^- \rightarrow O + 2e$	$5.47E-8T_e^{0.324}exp(-2.98/T_e)$
6	$e + O_2^+ \rightarrow 2O$	2.2E-8/Te
7	$O^- + O_2^+ \rightarrow O + O_2$	2.6E-8(300/T) ^{0.44}
8	$O^- + O \rightarrow O_2 + e$	(1.9, 3, 5)E-10
9	$O^- + O_2^+ \rightarrow 3O$	2.6E-8(300/T) ^{0.44}
Addition of C) ⁺	
10	$e + O_2 \rightarrow O^- + O^+ + e$	$7.1E - 11T_e^{0.5}exp(-17/T_e)$ $1.88E - 10T_e^{1.699}exp(-16.81/T_e)$
11	$e + O_2 \rightarrow O + O^+ + 2e$	$1.88E - 10T_{e}^{1.699}exp(-16.81/T_{e})$
12	$e + O \rightarrow O^+ + 2e$	9.0E-9Te ^{0.7} exp(-13.6/Te)
13	$O^- + O^+ \rightarrow 2O$	4.0E-8(300/T) ^{0.44}
14	$O^+ + O_2 \rightarrow O + O_2^+$	2.0E-11(300/T) ^{0.5}
Addition of n	netastable $O_2^{*}({}^1\Delta_v)$; see note f below	
15	$e + O_2 \rightarrow O_2^* + e$	1.37E-9 exp(-2.14/Te)
16	$e + O_2^* \rightarrow e + O_2$	2.06E-9 exp(-1.163/Te)
17	$e + O_2^* \rightarrow O + O^-$	$4.19E - 9T_e^{-1.376} \exp(-5.19/T_e)$
18	$O_2^* + O_2 \rightarrow 2O_2$	2.2E-18(T/300)0.8
19	$O_2^* + O \rightarrow O_2 + O$	(1.0, 7)E-16
20	$O^- + O_2^* \rightarrow O_1 + e$	2.2E-11
21	$O^- + O_2^* \rightarrow O_2^- + O$	1.1E-11
Addition of n	netastable O(¹ D)	
22	$e + O_2 \rightarrow O + O^* + e$	3.49E-8 exp(-5.92/Te)
23	$e + O \rightarrow O^* + e$	4.54E-9 exp(-2.36/Te)
24	$e + O^* \rightarrow e + O$	8.17E-9 exp(-0.4/Te)
25	$e + O^* \rightarrow O^+ + 2e$	$9.0E - 9T_{e}^{0.7} \exp(-11.6/T_{e})$
26	$O^* + O \rightarrow 2O$	8.0E-12
27	$O^* + O_2 \rightarrow O + O_2$	(6.4, 7.0)E-12 exp(67/T)
28	$O^* + O_2 \rightarrow O + O_2^*$	1.0E-12
Addition of s	elected reactions for O_2^- and O_3	
29	$O^- + O_2 \rightarrow O_3 + e$	5E-15
30	$e + O_3 \rightarrow O_2^- + O$	1E-9
31	$e + O_3 \rightarrow O^- + O_2$	2.12E-9Te ^{-1.058} exp(-0.93/Te)
32	$O_2^- + O_2^+ \rightarrow 2O_2$	2E-7(300/T)0.5
33	$O_2^- + O^+ \rightarrow O_2 + O$	(1, 2)E-7(300/T) ^{0.5}
34	$O_3 + O_2 \rightarrow O_2 + O + O_2$	7.3E-10exp(-11400/T)
35	$O_3 + O \rightarrow 2O_2$	1.8E-11exp(-2300/T)

Complex reaction schemes for O_2 plasma - 3^{rd} order reactions

Number	Reaction	Rate Constant (cm ⁶ /s)		
Reactions among e, O_2 , O_2^+ , and O^-				
1	$e + e + O_2^+ \rightarrow e + O_2$	$1E - 19(0.026/T_e)^{4.5}$		
2	$e + O_2^+ + O_2 \rightarrow O_2 + O_2$	6E-27(0.026/Te)1.5, 1E-26		
3	$e + O + O_2 \rightarrow O^- + O_2$	1E-31		
4	$\mathrm{O}^- + \mathrm{O}_2^+ + \mathrm{O}_2 \rightarrow \mathrm{O} + \mathrm{O}_2 + \mathrm{O}_2$	2E-25(300/T) ^{2.5}		
5	$\mathrm{O} + \mathrm{O} + \mathrm{O}_2 \rightarrow \mathrm{O}_2 + \mathrm{O}_2$	$2.45E - 31T^{-0.63}$		
		1.3E - 32(300/T)exp(-170/T)		
6	$\mathrm{O} + \mathrm{O} + \mathrm{O} \rightarrow \mathrm{O}_2 + \mathrm{O}$	6.2E - 32exp(-750/T)		
Addition of O ⁺				
7	$e + e + O^+ \rightarrow e + O$	$1E - 19(0.026/T_e)^{4.5}$		
8	$e + O^+ + O_2 \rightarrow O + O_2$	6E-27(0.026/Te) ^{1.5} , 1E-26		
9	$O^- + O^+ + O_2 \rightarrow O_2 + O_2$	2E-25(300/T) ^{2.5} , 2E-25		
10	$O^- + O^+ + M \rightarrow O + O + M$	$2E-25(300/T)^{2.5}$		
11	$\mathrm{O^+} + \mathrm{O} + \mathrm{O_2} \rightarrow \mathrm{O_2^+} + \mathrm{O_2}$	1E-29		
Addition of metastable $O(^{1}D)$				
12	$\mathrm{O} + \mathrm{O}^* + \mathrm{O}_2 \rightarrow \mathrm{O}_2 + \mathrm{O}_2$	9.9E-33		
Addition of selected reactions for metastable $O_2^{*}(^{1}\Delta_{v})$, O_2^{-} , and O_3				
13	$e + O_2 + O_2 \rightarrow O_2^- + O_2$	1.4E-29(0.026/Te)		
		$\times \exp(100/T - 0.061/T_{e})$		
14	$e + O_2 + O \rightarrow O_2^- + O$	1E-31		
15	$O^- + O_2^+ + O_2 \rightarrow O_3 + O_2$	2E-25(300/T) ^{2.5}		
16	$O + O_2 + O_2 \rightarrow O_3 + O_2$	6.9E-34(300/T) ^{1.25} ,		
		6.4E-35 exp(663/T)		
17	$O + O_2 + O \rightarrow O_3 + O$	2.15E-34 exp(345/T)		
18	$\mathrm{e} + \mathrm{O}_2^* + \mathrm{O}_2 \rightarrow \mathrm{O}_2^- + \mathrm{O}_2$	1.9E-30		
19	$\mathrm{e} + \mathrm{O}_2^* + \mathrm{O} \rightarrow \mathrm{O}_2^- + \mathrm{O}$	1E-31		
20	$\mathrm{O}_2^- + \mathrm{O}^+ + \mathrm{M} \to \mathrm{O}_3 + \mathrm{M}$	$2E-25(300/T)^{2.5}$		
21	$\mathrm{O}_2^- + \mathrm{O}_2^+ + \mathrm{O}_2 \rightarrow \mathrm{O}_2 - \mathrm{O}_2 + \mathrm{O}_2$	$2E - 25(300/T)^{2.5}$		