## Plasma and Dry Micro/Nanotechnologies 4. Deposition

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### 4.1 Introduction to Deposition

Deposition may be considered as six sequential substeps:

- 1. Adsorption of arriving atoms and molecules on the surface.
- 2. Diffusion before becoming incorporated into the film.
- 3. **Reaction** of adsorbed species with each other and the surface to form the bonds of the film material.
- 4. Initial aggregation of the film material nucleation.
- 5. Development of a structure (**morphology**) as the film grows thicker. It includes topography (roughness) and crystallography.
- 6. Diffusional interactions within the bulk of the film and with the substrate



### Surface Reactions $\Rightarrow$ Deposition

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4.2 Adsorption

Consider a molecule approaching a surface from the vapor phase.

- A few atomic distances from the surface molecule begins to feel an attraction interaction with the surface molecules by van der Waals forces (London dispersion forces):
  - molecules/atoms without dipole moments (symmetrical or intert) interact due to oscillating dipoles, i.e. induced-dipole interaction
  - polar molecules (permanent dipoles) interact more strongly
- The approaching molecule is being attracted into a potential well accelerates down the curve until it passes the bottom and is repelled by steeply rising potential.



An illustrative Lennard-Jones potential model for physical adsorption.

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Physisorption			

- If enough of the molecule's perpendicular component of momentum is dissipated into the surface the molecule cannot escape the well after being repelled 
  physisorption
  - fraction of physisorbed molecules trapping probability δ
  - reflected 1 − δ
  - $\delta$  is different from thermal accomodation coefficient  $\gamma$  introduced previously
  - molecule is at least partially accomodated thermally to the surface temperature T<sub>s</sub> even when it is reflected
- The physisorbed molecule is mobile on the surface except at cryogenic T hopping (diffusing) between surface atomic sites.



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Chemisorption			

During surface diffusion the molecule

- may desorb after a while by gaining enough energy in the tail of the thermal energy distribution.
- may undergo a further interaction consisting of the formation of chemical bonds with the surface atoms, i. e. chemisorption. The chemisorption reaction probability  $\xi$  is used in the case of chemisorption on a foreign substrate instead of condensation coefficient  $\alpha_c$ .
- Some of adsorbed species eventually escape back into the vapor phase ⇒ sticking coefficient S<sub>c</sub> fraction of the arriving vapor that remains adsorbed for the duration of the experiment.



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Chemisorption			

- Sticking coefficient S<sub>c</sub> has less fundamental meaning than δ and ξ (or α<sub>c</sub>) that are determined solely by chemistry and energy. Nevertheless, S<sub>c</sub> is very useful in thin film deposition it is the fraction of arriving vapor incorporated into the film (buried before escaping).
- Utilization fraction  $\eta$  of a chemical vapor fraction of molecules utilized for the deposition  $\Rightarrow \eta$  can approach unity even when  $S_c$  is very low.



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Precursor State			

Chemisorption - sharing electrons, physisorption - dipole interaction. If both adsorption states exist, the physisorption is called **precursor state**. The **precursor model** may also be applied to cases where both the adsorption states involve chemical bonding but one is weaker.

The precursor model has long been applied to heterogeneous catalysis, thin-film deposition and condensation of molecular vapors. Recent theory indicates that even the condensation of monoatomic vapor such as AI can involve both adsorption states: the precursor state = AI-AI dimer with limited bonding to bulk AI.

In most CVD reactions, the feed vapors adsorb as molecules that undergo reaction breaking their molecular bonds and form new bonds to surface

$$\operatorname{SiH}_4(g) \to \cdots \to \operatorname{SiH}_4(a) \to \operatorname{Si}(c) + 2\operatorname{H}_2(g)$$
 (1)

In deposition of compounds from separate vapor sources of each element, adsorbing vapor bonds much more strongly to surface sites occupied by the other element

$$\operatorname{Zn}(g) + \operatorname{Se}(a) \to \operatorname{ZnSe}(c) \quad \text{and} \quad \operatorname{Se}(g) + \operatorname{Zn}(a) \to \operatorname{ZnSe}(c)$$
 (2)

- Si chemically passivated by H reacts with adsorbates mainly at those few sites that are missing H atom. On the H-passivated sites, adsorbates remain only physisorbed.
- Atomically flat surfaces often bond more strongly with adsorbates at atomic steps.



Consider hypothetical diatomic gas-phase molecules  $Y_2(g)$  adsorbing and then dissociative chemisorbing as two Y atoms:



Lifting **atomic Y** out of its potential well along **curve c** results in much higher molar potential energy  $E_p$  in the gas phase - roughly the **heat of formation**,  $\Delta_f H$ , of 2*Y*(g) from  $Y_2(g)$ .

The **curve a** represents **activated chemisorption** - there is an activation energy  $E_{\rm a}$  to be overcome for  $Y_2(g)$  to become dissociatively chemisorbed.

For deeper precursor well, **curve b**, chemisorption is not activated though there is still a barrier  $E_{\rm rb}$ .

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## Eley-Ridel and Langmuir-Hinshelwood mechanisms

- In curve c, E<sub>p</sub> is high enough ⇒ direct chemisorption (without involving precursor state) ⇒ Eley-Rideal mechanism, i.e. direct reaction between an incoming species and a surface site
- Langmuir-Hinshelwood mechanisms: (i) adsorption from the gas phase, (ii) dissociation of molecules at the surface (iii) reaction among surface species (chemisorbed molecules)



Two ways in which vapor can arrive at surface having  $E_{\rm p}>$  0:

- Gaseous molecules have their E<sub>p</sub> raised by becoming dissociated.
- Solids and liquids have it raised by evaporating.

Energy-enhanced deposition processes provide enough energy  $E_{\rm p} > E_{\rm a}$ 

- sputter deposition arriving species have kinetic energy ~ 1000 kJ/mol and E<sub>p</sub> > 0.
- plasma-enhanced deposition vapor molecules are dissociated in plasma

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## Thermally controlled deposition process

In evaporation and CVD, the vapors often adsorb first into the precursor state (curve a, b)

- $\blacktriangleright$  and chemisorb by overcoming the barrier  $\textit{E}_{r(a,b)}$
- or desorb by overcoming the heat of physisorption  $E_{d(a,b)}$ .

 $\Rightarrow$  these two reactions result in a net rate of chemisorption.

### First-order chemical reaction

$$R_k = k_k n_s = k_k n_{s0} \Theta \tag{3}$$

 $R_k$  rate of kth surface reaction per unit area  $[m^{-2}s^{-1}]$  $k_k$  rate constant  $[s^{-1}]$  $n_s$  surface concentration of reactant  $[m^{-2}]$  $\Theta$  fractional surface coverage by reactant with rate constant following Arrhenius equation:

$$k_k = \nu_{0k} \exp(-\frac{E_k}{RT}) \tag{4}$$

 $\nu_{0k}$  frequency factor or pre-exponential factor

 $E_k$  reaction activation energy [kJ/mol]

**mass balance for the physisorbed precursor**: assuming  $n_s$  is constant over time (steady state), chemisorption reaction occurs only in the forward direction (not valid for too high T at which film begins to decompose), adsorption occurs on free sites

$$\Gamma_{i}\delta(1-\Theta) = R_r + R_d = (k_r + k_d)n_{s0}\Theta$$
(5)

 $\Gamma_i$  molecular impingement flux  $[m^{-2}s^{-1}]$ ,  $\delta$  trapping probability  $R_r$  reaction (chemisorption) rate and  $R_d$  desorption rate

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Sticking Coefficient			

Using previous relations we obtain for surface coverage  $\Theta$ 

$$\Theta = \frac{\Gamma_{\rm i} \delta / n_{\rm s0}}{\Gamma_{\rm i} \delta / n_{\rm s0} + k_r + k_d} \tag{6}$$

and substituing into the chemisorption rate expression (3)

$$R_r = k_r n_{s0} \Theta = \frac{\Gamma_i \delta k_r}{\Gamma_i \delta / n_{s0} + k_r + k_d}.$$
(7)

We may now define sticking coefficient S<sub>c</sub> more precisely

$$S_c = R_r / \Gamma_i$$
 (8)

In case of small  $\Gamma_i$ , i.e. small  $\Theta$  Eq. (7) simplifies to

$$R_{r} = \frac{\Gamma_{i}\delta}{1 + k_{d}/k_{r}} = \Gamma_{i}\frac{\delta}{1 + \frac{\nu_{0d}}{\nu_{0r}}\exp(\frac{E_{r} - E_{d}}{RT_{s}})} \equiv \Gamma_{i}\xi.$$
(9)

where  $\xi$  is chemisorption reaction probability.

 $S_c \approx \xi$  for  $\Theta \ll 1$  but for larger  $\Theta$ ,  $S_c < \xi$ .

For the special case of film deposition from single vapor having the same composition as the film,  $\xi = \alpha_c$  (condensation coefficient).

The assumption of 1st order kinetics is not always valid for a more complicated case of compound-film deposition from multicomponent vapors. For more details refer to book of Donald Smith, chapter 7.3.3.

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### Activated adsorption

The chemisorption rate  $R_r$  governs the rate of film deposition when

- $k_r$  is the same from site to site along the surface
- $rac{T_s}$  is not so high that decomposition or re-evaporation of the film occur.



$$R_r = \Gamma_i \frac{\delta}{1 + \frac{\nu_{0d}}{\nu_{0r}} \exp(\frac{E_r - E_d}{RT_s})}$$

If  $E_r > E_d$  there is an **activation energy**   $E_a = E_{ra} - E_{da}$  for chemisorption (**curve a**). For high  $E_a$  the film fail to deposit unless  $T_s$  is raised to make exp. term smaller  $\Rightarrow R_r \uparrow$  for  $T \uparrow$ 

The activated case is very common in CVD.

If  $E_r < E_d$  (chemisorption is not activated as in **curve b**)  $\Rightarrow R_r \downarrow$  for  $T \uparrow$ 

**Example**: Si deposits from SiH<sub>4</sub> at elevated (not room)  $T_s$ . If  $T_s$  becomes too high, the Si evaporation flux  $\Gamma_v > R_r$  and deposition stops. The net deposition flux of Si:

 $\Gamma_r = R_r - \Gamma_v$  (10)  $\Rightarrow T_s$  window for deposition.

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## 4.3 Surface Diffusion

typical E<sub>p</sub>, k.J/mol E. -400 site 1 site 2 transition state -600 physisorption chemisorption transition state site 1 site 2 (b)

**Flux of adsorbate**  $\Gamma_s$  [m<sup>-1</sup>s<sup>-1</sup>] across the  $E_s$  barrier between 1 and 2 sites in the *x* direction

#### Rate of barrier crossing R<sub>s</sub>

$$\mathbf{R}_s = \frac{\mathbf{\Gamma}_s}{a}$$
 (11)

(a is the distance between the sites)

Considering Maxwell-Boltzmann distribution

$$R_s = \frac{1}{4} \frac{n'_s \bar{v}}{a} = \frac{n'_s}{a} \sqrt{\frac{k_{\rm B} T}{2\pi m}} \quad (12)$$

 $n'_{s}$  - surface concentration of adsorbate residing in the transition state.

Relation between  $n'_s$  and  $n_s$ ? - from statistical physics (next page)  $\Rightarrow$  Arrhenius law for  $R_s$ 

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### **Concentrations in Surface States**

Relation between  $n'_s$  (concentration in transition states) and  $n_s$  (concentration in adsorption states)?

Z - partition function,  $g_i$  degeneracy of the energy level,

r, v, t rotational, vibrational, translation kinetic energies (electronic excitations are neglected at ordinary T),

The Boltzmann factor accounts for the potential energy difference  $E_s$  (J/mol) between the adsorption-site state  $n_s$  and the transition state  $n'_s$ 

- rotation of molecules is limited by adsorption  $\Rightarrow Z'_r = Z_r = 1$
- Z<sub>vk</sub> derived from quantum mechanics for harmonic oscillator

$$Z_{\nu k} = \frac{1}{1 - \exp\left(-\frac{h\nu_k}{k_{\rm B}T}\right)} \tag{14}$$

but vibrational modes are mostly in their ground states at ordinary T

partition function for translation energy

$$Z_{tx}' = a \frac{\sqrt{2\pi m k_{\rm B} T}}{h} \tag{15}$$

$$\Rightarrow n'_{s} = n_{s} Z'_{tx} \exp \left(-\left(\frac{E_{s}}{RT}\right)\right)$$
(16)

## Rate of barrier crossing $R_s$ , molecular hopping rate $k_s$

... using absolute-reaction-rate theory (predicts the absolute reaction rate of a chemical reaction from the quantum mechanical description of the potential energy changes during the interaction; cannot provide a quantitative estimation of the diffusion rate but gives an insight into the determining factors)

$$R_{s} = n_{s} \left(\frac{k_{\rm B}T}{h}\right) \exp\left(-\frac{E_{s}}{RT}\right) = n_{s}\nu_{0s} \exp\left(-\frac{E_{s}}{RT}\right) = n_{s}k_{s}$$
(17)

Arhenius expression for the rate constant  $k_s$  of "chemical reaction".

The rate constant  $k_s$  (s<sup>-1</sup>) represents the frequency with which an individual adsorbate molecules "hops" to an adjacent site. Thus, the factor  $\nu_{0s} = 10^{13} - 10^{16}$  s<sup>-1</sup> is NOT the frequency of any vibrational component  $\nu_k$  of the adsorbate.

The rate of surface diffusion increases exponentially with  $\uparrow~{\cal T}$  and  $\downarrow~{\cal E}_{\rm s}$  (activation energy for surface diffusion).

 $E_s \ll E_d, E_c$  (desorption activation energy of physisorbed or chemisorbed species, respectively)

 $\Rightarrow$  high rate of surface diffusion at film deposition *T* approaching the onset of re-evaporation, i.e. when  $\exp(-\frac{E_c}{RT})$  becoming significant = one of principal ways in which substrate *T* affects film structure.

 $E_s/E_c$  sometimes referred as corrugation ratio, it is lower for metals than for semiconductors due to absence of bond directionality in metals.

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Diffusion Length			

Relation between **molecular hopping rate**  $k_s$  and the distance which an adsorbate molecule travels during the film deposition:

- classic random-walk problem
- ► The final locations are more widely dispersed from the starting point with increasing time *t*. For large number of hops  $N_o$ , it is Gaussian dispersion  $\exp(-\frac{x^2}{2\sigma^2})$  characterized by its standard deviation  $\sigma$  here the diffusion length  $\Lambda$

$$\Lambda = r\sqrt{N_o} \approx a\sqrt{N_o} = a\sqrt{k_s t}$$

r is per-hop rms change in the distance from the starting points

How  $\land$  depends on temperature? Two regimes need to be considered separately:

- t is time between adsorption and burial by the next depositing monolayer
- adsorbate is more likely to desorb than to be buried within t

Regime 1 - burial time:

$$t = rac{n_0}{\Gamma_{
m r}} \quad \Rightarrow \quad \Lambda = a \sqrt{rac{
u_{
m os} n_0}{\Gamma_{
m r}}} \exp{\left(-rac{E_{
m s}}{2RT}
ight)}$$

where  $n_0$  is surface density of adsorption sites (m<sup>-2</sup>),  $\Gamma_r$  deposition flux (m<sup>-2</sup>s<sup>-1</sup>) and  $k_s = \nu_{os} \exp\left(-\frac{E_s}{RT}\right)$ 



## Diffusion Length - contin.

*Regime 2* - If *T* is high enough the film re-evaporation (**desorption**) becomes significant.  $\Rightarrow$  *t* is adsorption lifetime. Considering the desoprtion only from chemisorbed state (negligible concentration of precursor state at high *T*):

$$t = rac{1}{k_{
m c}} = rac{1}{
u_{
m oc}} \exp{(rac{E_{
m c}}{RT})}$$

where subscript c denotes the chemisorbed state.

Combining both gives

$$\Lambda = a \sqrt{\frac{\nu_{\rm os}}{\nu_{\rm oc}}} \exp\left(\frac{E_{\rm c} - E_{\rm s}}{2RT}\right) \tag{18}$$



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<b>Diffusion Coefficient</b>			

How is the diffusion length  $\land$  (obtained from the examination of motion of individual adsorbed molecules) related to the macroscopic quantity - diffusion coefficient? Transport equation

$$\Gamma_{\rm s} = -D \frac{{\rm d}n_{\rm s}}{{\rm d}x} \qquad \qquad \Gamma_{\rm s}, \, n_{\rm s} \text{ have surface units } m^{-1} s^{-1} \text{ and } m^{-2}$$

Using analogy to 3D case  $D = \frac{1}{4} v_{av} \lambda$ , in which the mean free path  $\lambda$  is the hop distance *a*, and the mean speed  $v_{av}$  is  $k_s a$ 

$$D = \frac{1}{4}k_{\rm s}a^2 \quad \Rightarrow \quad \Lambda = a\sqrt{k_{\rm s}t} = 2\sqrt{Dt} \tag{19}$$

Thus, we can express D in the Arrhenius form when using this expression for  $k_s$ :

$$D = \frac{1}{4}\nu_{\rm os}a^2 \exp\left(-\frac{E_{\rm s}}{RT}\right) = D_0 \exp\left(-\frac{E_{\rm s}}{RT}\right)$$
(20)

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4.4 Nucleation			

Nucleation is a complication that must often be added to the above described model of chemisorption where we expected that  $k_r$  is the same from site to site.

If nucleation is important, the net deposition flux  $\Gamma_r = R_r - \Gamma_v$  is  $\Gamma_r > 0$  only for certain active substrate-surface sites, **nucleation sites** or **nuclei of film material** which have spontaneously accumulated. Various examples:

► Si chemically passivated by H reacts with adsorbates mainly at those few sites that are missing H atom  $\Rightarrow$   $R_r \uparrow$  at unpassivated Si surface atoms because of  $\downarrow E_r$ 

In deposition of compounds from separate vapor sources of each element, adsorbing vapor bonds much more strongly to surface sites occupied by the other element

 $Zn(g) + Se(a) \rightarrow ZnSe(c)$  and  $Se(g) + Zn(a) \rightarrow ZnSe(c)$ 

► atoms of low-reactivity metals often bond much less easily to nonmetallic substrates: *E*<sub>a</sub> of chemical bond of Zn, Cd to glass is very high because of high bond strength between substrate elements ⇒ Zn, Cd on glass bonds more readily to itself than to the surface ⇒ formation of nucleus

The existence of certain sites which are active in adsorption is common in thin-film processes. Access of precursor to these favored sites can dominate the deposition kinetics. Two types of access

- from vapor phase Eley-Rideal mechanism
- by surface diffusion Langmuir-Hinshelwood mechanisms

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### Surface tension $\gamma$ and surface free energy

Concept of surface tension  $\gamma$  has to be introduced to understand nucleation.

The force *F* required to draw a liquid membrane:  $F = 2b\gamma$  (*b* is circumference, 2 stands for inner and outer surface). Work  $F\Delta x$  to create the membrane of area  $A = 2b\Delta x$  is stored as surface energy  $\Rightarrow$  surface energy per unit area  $F\Delta x/A = \gamma$  [N/m]

 $\Rightarrow$  For liquids, surface (free) energy per unit area (J/m<sup>2</sup>) is equal to surface tension  $\gamma$ .



**For solids at** T > 0 **K**, the surface Gibbs free energy is reduced by entropy factor G = H - TS which depends on the degree of surface disorder  $\Rightarrow$  surf. energy is minimized by surface diffusion.

In solids, there is an surface energy term  $A \sum_{ij} \sigma_{ij} d\varepsilon_{ij}$  in which  $\sigma_{ij}$  is surface stress and  $d\varepsilon_{ij}$  surface strain tensor  $\sum d\varepsilon_{ij} \delta_{ij} = dA/A$ . Liquids cannot support such strain because the atoms just rearrange to relax it.

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Thin-film growth			

Nature tends to minimize surface energy  $\gamma A$ :

- when wire is lifted far enough, the membrane is in the plane of the ring
- ▶ in solids, surface energy minimizes by surface diffusion ⇒ fundamental process to development of structure in thin films

#### In thin-film growth, both A and $\gamma$ are varying:

- A depends on surface topography
- γ depends on many properties of exposed surface (chemical composition, crystallographic orientation, atomic reconstruction, atomic-scale roughness etc.). It is anisotropic for most crystalline solids.

For deposition onto a foreign substrate, nucleation behaviour is strongly influenced by surface tension of substrate,  $\gamma_s$ . We also need to consider  $\gamma_i$  of the substrate-film interface and  $\gamma_f$  of film free surface.

Two cases for growth modes:

- $\Lambda \gg a$ , i. e. deposition material can rearrange itself to minimize  $\gamma$ , nucleation is not kinetically limited and approaches equilibrium
- ► Λ < a atoms sticks where they land and the growth behaviour is "quenched"

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Growth Modes for $\Lambda \gg$	> a	(nucleation is not kinetic	ally limited)	
For $\Lambda \gg a$ , there are two nucleation situations on the bare substrate				

- (a) films wets the substrate because  $\gamma_f + \gamma_i < \gamma_s \Rightarrow$  smooth growth, atomic layer by layer (Frank-van der Merwe growth). It requires strong enough bonding between film and substrate to reduce  $\gamma_i$
- ▶ (b) with insufficient substrate bonding film forms 3D islands (Volmer-Weber growth mode). In extreme case of no bonding at all *γ*<sub>i</sub> = *γ*<sub>f</sub> + *γ*<sub>s</sub>, the film spreading across the substrate would increase the total surface energy by 2*γ*<sub>f</sub>.





Third growth mode, **Stranski-Krastanov**, shown in (c) - growth mode changes from layer to island after a monolayer or two due to a change in the energy situation.



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2D versus 3D nucleati	on		
3D nucleation is usually	undesirable since		

it leads to rough, nonuniform films (extreme example being the diamond nuclei)



How to manipulated with the growth mode? Film wetting for  $\gamma_f + \gamma_i < \gamma_s$ .

 $\gamma_{\rm i}$  decreases with film-substrate bonding (covalent, ionic, metallic). In general, interfacial bonding is stronger between materials having the same type of bonding. Examples:

- chemically-active metal, Cr, will bond to glass by breaking Si-O bond and forming Cr-O
- Au cannot do this, i.e. does not bond well to glass
- Au forms a strong metallic bond to clean Cr

 $\Rightarrow$  Using intermediate "glue" layer (Cr) which bonds well to both, the film and the substrate,  $\gamma_i$  can be reduced and wetting accomplished. Another good bonding material is Ti.

Alternatively, energy-enhanced techniques (plasma treatment, ion bombardment, sputtering) can provide the activation energy for bonding between film and substrate, i.e. reduce  $\gamma_i$ .

#### Two ways in which 3D nuclei can form

- bonding initiates at active surface sites such as atomic steps, crystal deffects, impurities. At these nucleation sites, the *E<sub>a</sub>* is lower than elsewhere.
- even if there are no active nucleation sites, 3D nuclei can still form at random surface locations because of the interfacial bonding which develops by the spontaneous accumulation of mobile atoms plus arriving vapor into "critical" nuclei which are big enough to be stable (classical nucleation problem)



## 3D nucleation - how to produce smooth & uniform film?

- ► concentration of critical nuclei n\* has to be ↑
- ▶ and their radius  $r^* \downarrow$ ,
- i.e., less coarse nucleation

How to achieve it?

- ► Using very high vapor arrival rate (supersaturation), at least until the nucleation phase is over and the film is continuous. Coarsening will still occur even with one-atom critical nuclei because atoms and nuclei are mobile on the surface ⇒ migration and coalescence
- ► Decrease the substrate *T* to inhibit surface diffusion ⇒ freezing the nucleation and coalescence. If arriving species do no thave enough energy to desorb or diffuse they remain where they land quenched growth mentioned earlier. In this case, the nucleation is kinetically inhibited by the surface-diffusion activation-energy barrier *E*<sub>s</sub>

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## 4.5 Structure Development

Upon coalescence of the surface nuclei to form a continuous film, the nucleation step of the film deposition is complete and **4th step begins - development of the bulk film structure**. The form of the structure changes dramatically with

- the amount of thermal motion taking place during film growth: scales with T<sub>s</sub>/T<sub>m</sub> (melting point of the film over substrate temperature - in K) - known as homologous or reduced T
- the amount of additional energy being delivered to the growth surface



Figure 5.15 Characteristics of the four basic structural zones and of whiskers, in cross section. The ratio of substrate T to film melting T  $(T_g/T_m)$  increases in the direction  $Z1 \rightarrow ZT \rightarrow Z2 \rightarrow Z3$ .

### 4.5 Structure Development

### quenched growth for Z1 and ZT:

- Z1 occurs at T<sub>s</sub>/T<sub>m</sub> so low that surface diffusion is negligible, i. e. A < a. Z1 consits of columns typically tens of nm in diameter separated by voids a few nm across. The columns have poor crystallinity (many deffects) or are amorphous. In thicker films, an array of cones with wider voids between them become superimposed upon this structure. The cones terminate in domes at the surface, and the size of the domes increases with film thickness.
- ZT also occurs when Λ < a. It contains defected columns as Z1 but the voids and domes are absent. ZT is usually associated with energy-enhanced processes.</p>



Figure 5.15 Characteristics of the four basic structural zones and of whiskers, in cross section. The ratio of substrate T to film melting T  $(T_s/T_m)$  increases in the direction  $Z1 \rightarrow ZT \rightarrow Z2 \rightarrow Z3$ .

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## 4.5 Structure Development

thermally activated rearrangement on or within the film for Z2 and Z3:

- ► Z2 occurs for  $T_s/T_m \gtrsim 0.3$ , diffusion becomes significant. It consits of columns having tight grain boundaries between them and having a characteristic diameter which increases with  $T_s/T_m$ . Crystalline columns are less defected than in Z1 and ZT, and are often facetted at the surface. The Z2 structure can also accur in amorphous films. The column boundaries are planes of reduced bonding rather than planes of crystallographic discontinuity.
- > Z3 occurs at certain instances at  $T_s/T_m > 0.5$ , considerable bulk annealing of the film is taking place during deposition, more isotropic crystallite shapes.



Figure 5.15 Characteristics of the four basic structural zones and of whiskers, in cross section. The ratio of substrate T to film melting T  $(T_s/T_m)$  increases in the direction Z1-2T-2Z-2Z3.



Simple model of statistical roughnening caused by statistical fluctuation in the vapor arrival flux: each atom is constrained to stick on the site it lands on (even if it is on top of a pillar)



Figure 5.16 Statistical roughening in random ballistic deposition of a 25-atom-thick film. (Pascal solution courtesy of Jared Smith-Mickelson.)

For large enough  $\bar{N}$  (average number of atoms per depositing site) the variation in heights is described by Gaussian distribution with standard deviation

$$\sigma = a\sqrt{ar{N}} = \sqrt{aar{h}}$$

where  $\bar{h} = a\bar{N}$  is average film thickness

Note: analogy to dispersion in lateral direction arising from surface diffusion because both are random processes

### **Quenched growth**



Atomistic processes in quenched-growth structure development (more realistic model):

- (a) ballistic aggregation (arriving atoms cannot perch on top of each other but rather settle sideways)
- (b) effect of atoms finite size (shadowing low areas)
- (c) sideways attraction (development of columns)
- (d) oblique shadowing (self-shadowing, incidence over range of θ occurs for fluid-flow regime Kn ≪ 1)
- (e) tilt effect
- (f) low sticking coefficient
- (g) void-filling by energetic particle due to enhanced mobility (left) and forward sputtering (right)

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## Quenched growth - void filling by energetic particles



2D MD simulation of the deposition of energetic atoms impinging perpendicularly onto a substrate held at 0 K.

The horizontal line is substrate interface.

 $E_t/E_c =$  incident energy / adatom potential-well depth

- (a)  $E_t/E_c = 0.02$
- ▶ (b) E<sub>t</sub>/E<sub>c</sub> = 0.5
- (c)  $E_t/E_c = 1.5$



## Thermally activated growth



(c) T = 440 K

(d) T = 460 K

Figure 5.20 Two-dimensional computer simulation of the effect of substrate T on void filling by surface diffusion. (Source: Reprinted from Ref. 22 by permission.)

## 4.6 Thin-Film Deposition Process Steps

All thin-film processes contain the four (or five) sequential steps.

1. A source of film material is provided.

Solid, liquid, vapor or gas source. Solid materials need to be vaporized (by heat or energetic beam of electrons, photons, i.e. laser ablation, or positive ions, i.e. sputtering) - **physical vapor deposition (PVD)**. The methods using gases, evaporating liquids or chemically gasified solids are **chemical vapor deposition (CVD)** methods.

2. The material is transported to the substrate.

The major issue is uniformity of arrival rate over the substrate area. Transport in a high vacuum = straight travelling lines  $\rightarrow$  importance of geometry. Transport in a (gaseous) fluid = many collisions  $\rightarrow$  gas flow patterns, diffusion of source molecules through other gases present.

3. The film is **deposited** onto the substrate surface.

It is influenced by source and transport factors and the conditions at the deposition surface. Three principal surface factors: (i) surface condition (roughness, contamination, degree of chemical bonding with the arriving materials and crystallographic parameters in the case of epitaxy), (ii) reactivity of arriving material (sticking coefficient  $S_c$  from 1 to less than  $10^{-3}$ ) and (iii) energy input (substrate heating, photons, ions, chemical energy).

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Thin-Film Deposition	Process Stens		

- 4. (Optionally, annealing takes place)
- 5. The final step is **analysis** of the film.

One level of analysis is the determination of functional properties important for given application and optimization of the deposition process for these processes (emphirical approach). A deeper level of analysis involves probing the film structure and composition (better understanding of the overall processes).

Analysis of the films after deposition - kind of final process monitoring. However, **monitoring** is important in all steps!

## 4.7 Overview of Deposition Methods I

method/processes	specification	
evaporative techniques:		
thermal (vacuum) evaporation	resistive heating	
	flash evaporation	
	arc evaporation	
	exploding-wire technique	
	rf heating	
	electron-beam evaporation	
pulsed laser deposition (PLD)		
molecular beam epitaxy (MBE)		
liquid-phase chemical techniques:		
electro processes	electroplating	
	electrolytic anodization	
mechanical techniques	spray pyrolysis	
liquid phase epitaxy		
gas-phase ch	emical techniques:	
chemical vapor deposition (CVD)	CVD epitaxy	
	metalorganic CVD (MOCVD)	
	low-pressure CVD (LPCVD)	
	atmospheric-pressure CVD (APCVD)	
atomic layer deposition (ALD)		
gas-phase physical-chemical techniques		
(except plasma and ion beam):		
modifications of CVD	hot filament CVD (HFCVD)	
	laser-induced CVD (PCVD)	
	photo-enhanced CVD (PHCVD)	
	electron enhanced CVD	

## **Overview of Deposition Methods I - evaporative methods**

#### vacuum evaporation



#### pulsed laser deposition



# vacuum evaporation (resistive and electron beam



#### molecular beam epitaxy



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## **Overview of Deposition Methods II**

method/processes	specification
ρ	lasma techniques:
sputter deposition	dc sputtering
	rf diode sputtering
	magnetron sputtering
PECVD in low temperature	dc discharge
discharges	rf capacitively coupled plasma (CCP)
	rf inductively coupled plasma (ICP)
	microwave ECR deposition
	microwave resonantor reactor
	atmospheric pressure dielectric barrier discharge (DBD)
	atmospheric pressure glow discharge (APGD)
	atmospheric pressure surface barrier discharge
	etc.
plasma processing in high temperature	vacuum arc
discharges	dc torch
	microwave torch
	etc.
ion beam techniques:	
sputter deposition	ion beam sputter-deposition
	reactive ion beam sputter-deposition
ion deposition	ion beam deposition
	ionized cluster beam deposition (ICB)
dual processes	ion beam assisted deposition (IBAD)
	dual ion beam deposition

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## **Overview of Deposition Methods II - ion beam**

### ion beam sputter-deposition



### ion-beam assisted deposition (IBAD)



### dual ion-beam deposition



