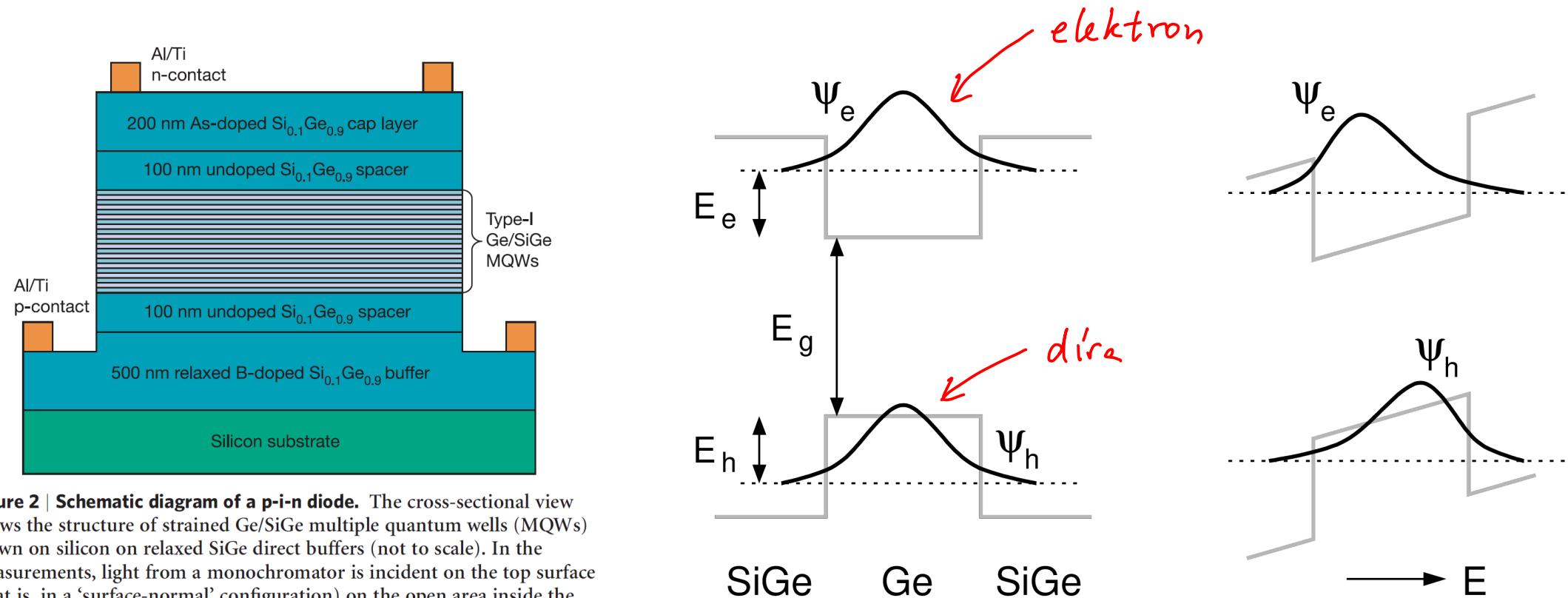


# Strong quantum-confined Stark effect in germanium quantum-well structures on silicon

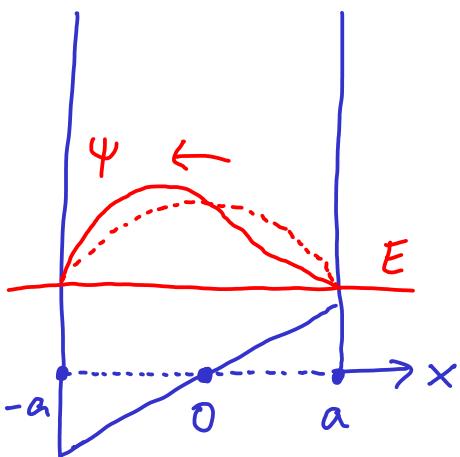
Yu-Hsuan Kuo<sup>1</sup>, Yong Kyu Lee<sup>1</sup>, Yangsi Ge<sup>1</sup>, Shen Ren<sup>1</sup>, Jonathan E. Roth<sup>1</sup>, Theodore I. Kamins<sup>1,2</sup>, David A. B. Miller<sup>1</sup> & James S. Harris<sup>1</sup>



**Figure 2 | Schematic diagram of a p-i-n diode.** The cross-sectional view shows the structure of strained Ge/SiGe multiple quantum wells (MQWs) grown on silicon on relaxed SiGe direct buffers (not to scale). In the measurements, light from a monochromator is incident on the top surface (that is, in a ‘surface-normal’ configuration) on the open area inside the rectangular frame top electrode (that is, between the portions of the AlTi n-contacts shown in this cross-section).

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# Kvantová jáma v elektrickém poli



→  $\bar{E}$

- 1) zkušební vlnová funkce  $\Psi(x) = N \left(1 + \lambda \frac{\pi x}{a}\right) \cos \frac{\pi x}{2a}$
- 2) normování
- ↑  
posouvá vpravo pro  $\lambda > 0$
- ↑  
 $\Psi$  bez pole

$$\int_{-a}^{+a} |\Psi(x)|^2 dx = 1 \rightarrow N$$

- 3) střední hodnota  $\hat{H}$

$$\langle H \rangle = \int_{-a}^{+a} \Psi^*(x) \left( -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + e \mathcal{E} x \right) \Psi(x) dx$$

$$\approx \frac{\hbar^2}{2ma^2} \left( \frac{\pi^2}{4} + \pi^2 \lambda^2 \right) + e \mathcal{E} a \frac{2(\pi^2 - 6)}{3\pi} \lambda$$

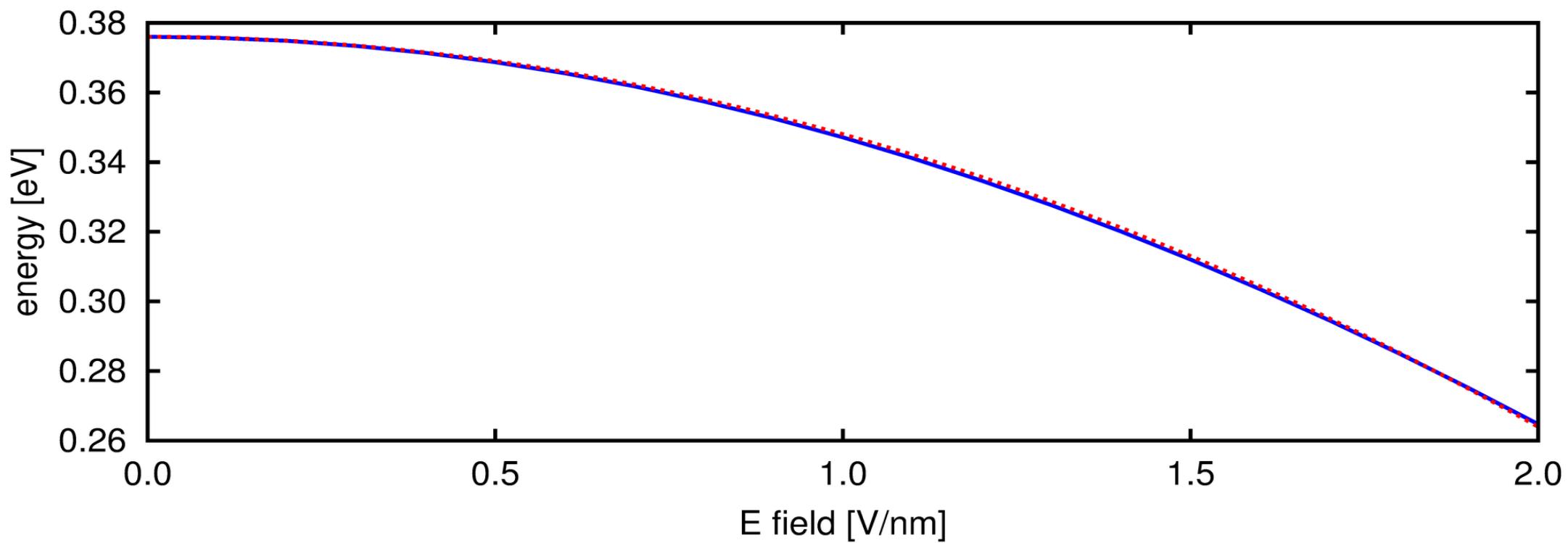
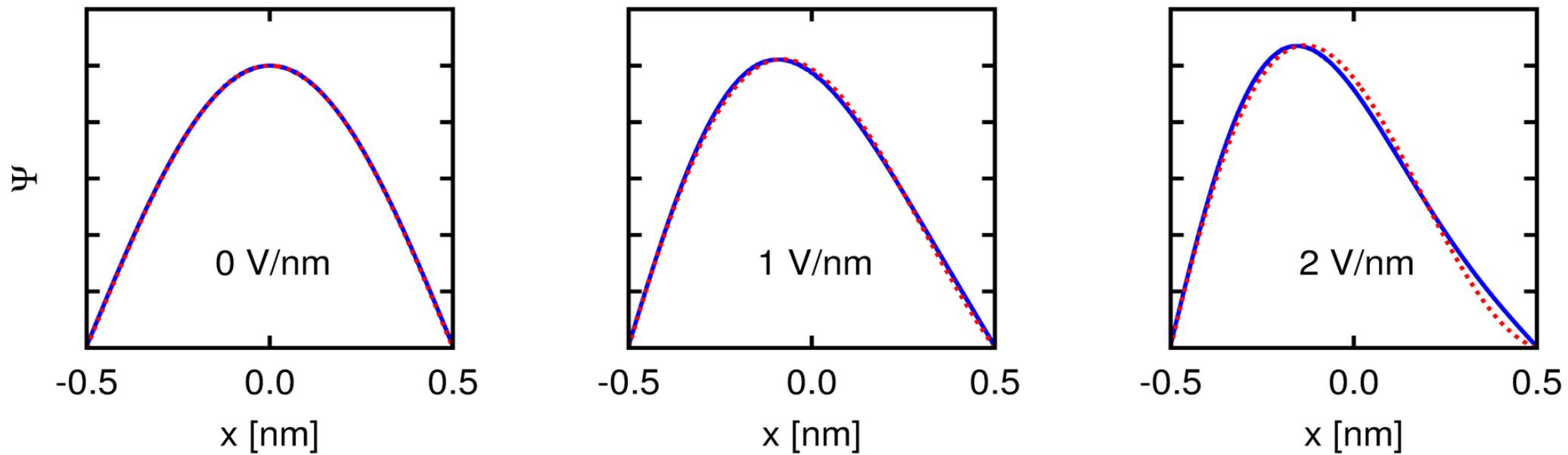
- 4) minimalizace

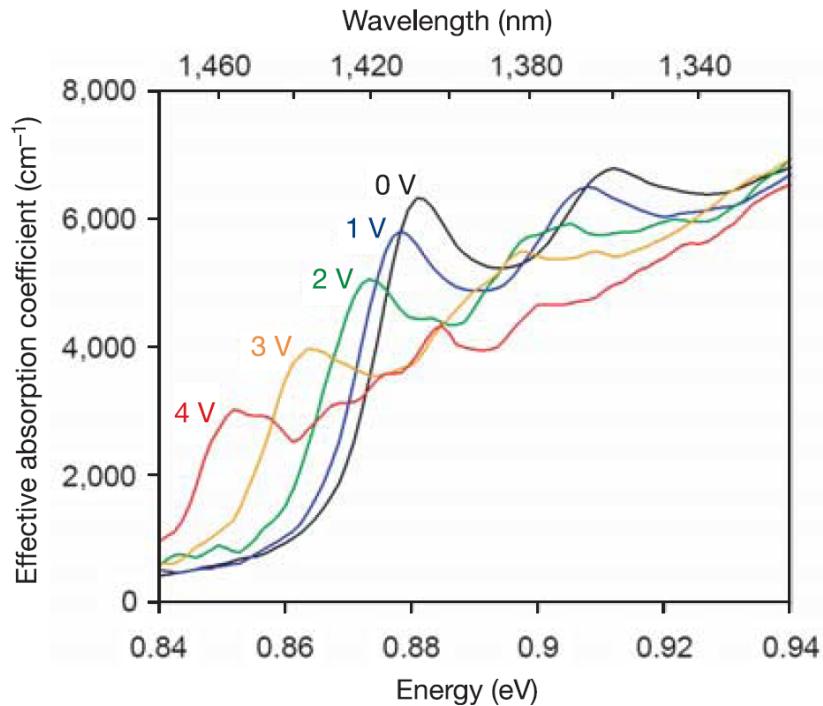
$$\frac{\partial \langle H \rangle}{\partial \lambda} = 0 \rightarrow \lambda \approx - \frac{\pi^2 - 6}{3\pi^3} \frac{e \mathcal{E} a}{\frac{\hbar^2}{2ma^2}}$$

posuv  $\sim \mathcal{E}$

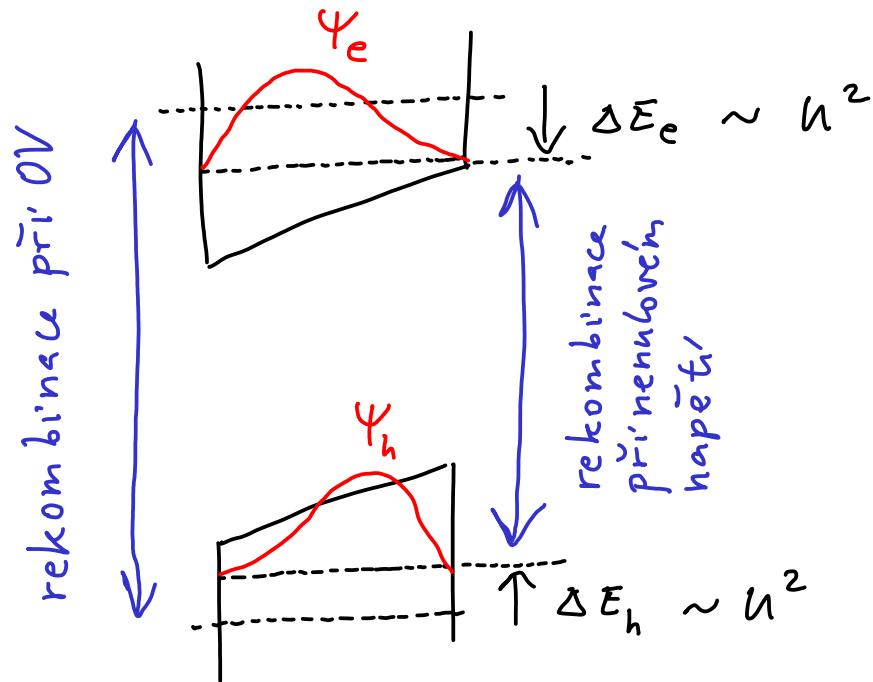
$$\langle H \rangle_{\text{opt}} \approx E_0 - \left( \frac{1}{3} - \frac{2}{\pi^2} \right)^2 \frac{(e \mathcal{E} a)^2}{\frac{\hbar^2}{2ma^2}}$$

snižem energii  
 $\sim \mathcal{E}^2$





**Figure 3 | Effective absorption coefficient spectra.** Strong QCSE is observed at room temperature with reverse bias from zero to 4 V. The thickness for effective absorption coefficient calculations is based on the combination of Ge well and SiGe barrier thicknesses.



1) poloha absorpčního píku

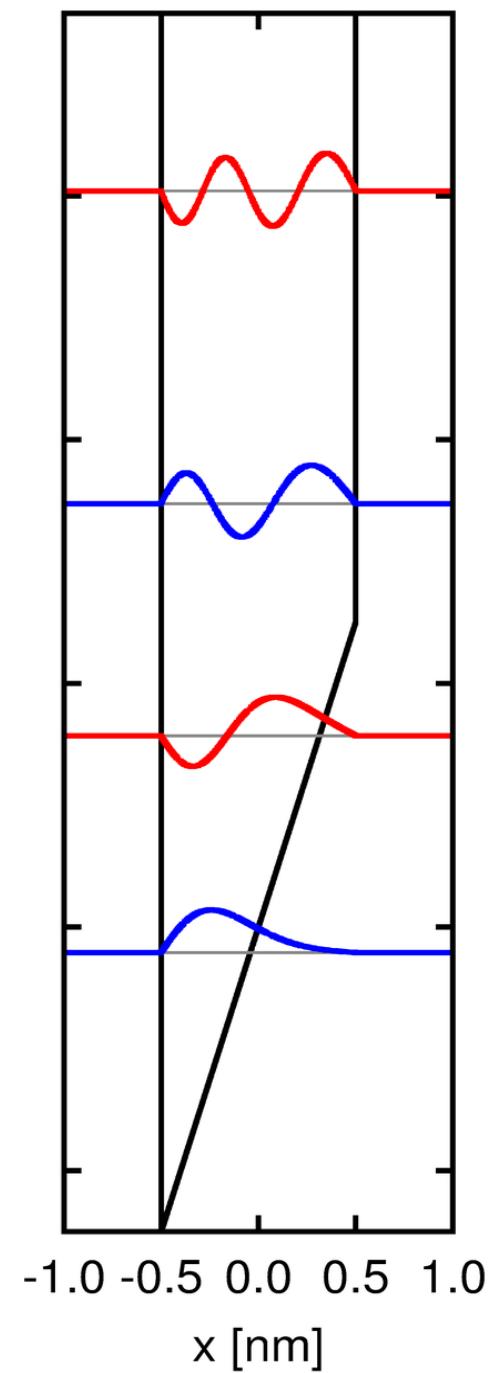
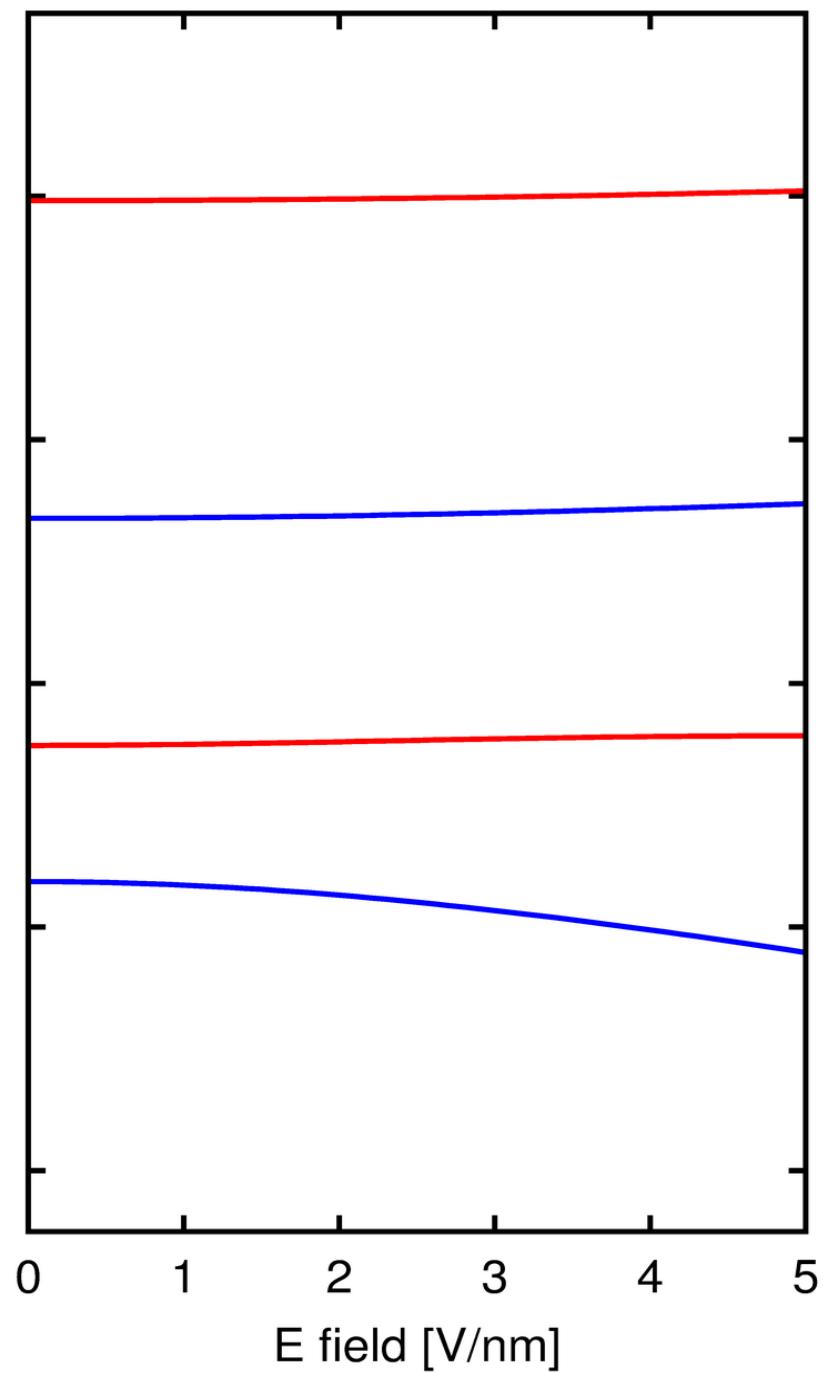
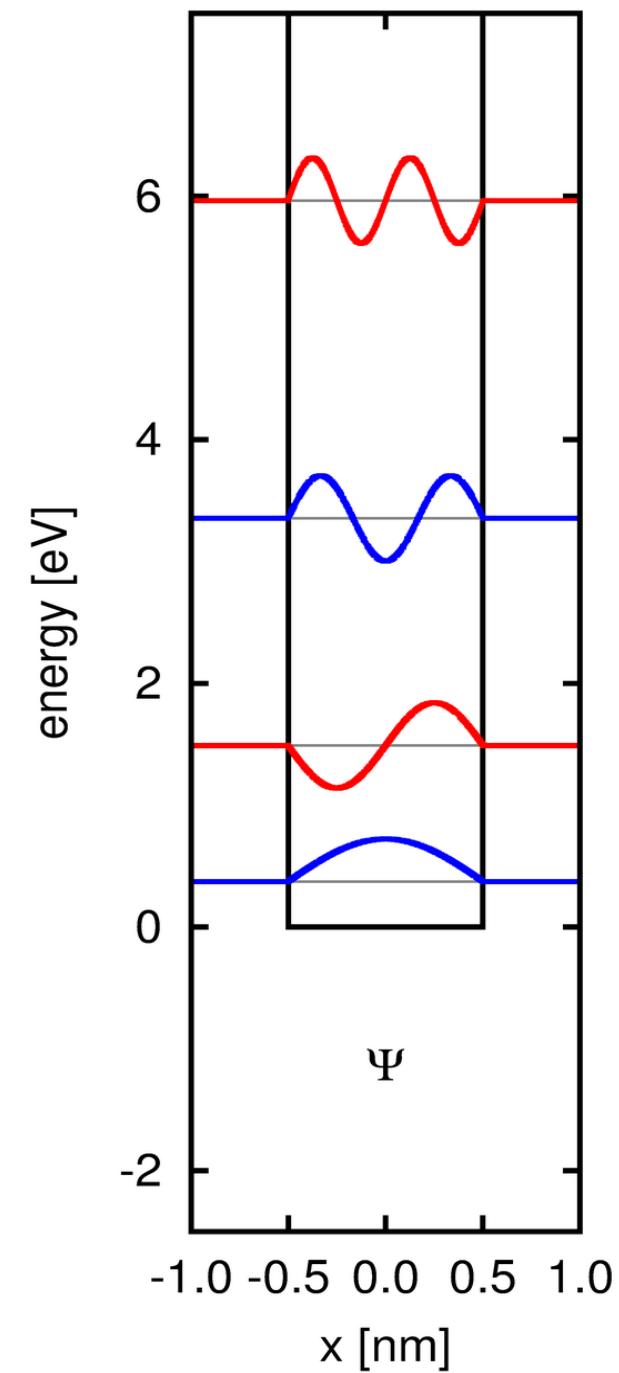
$$E_{rec} = E_{rec}(0V) - \alpha u^2$$

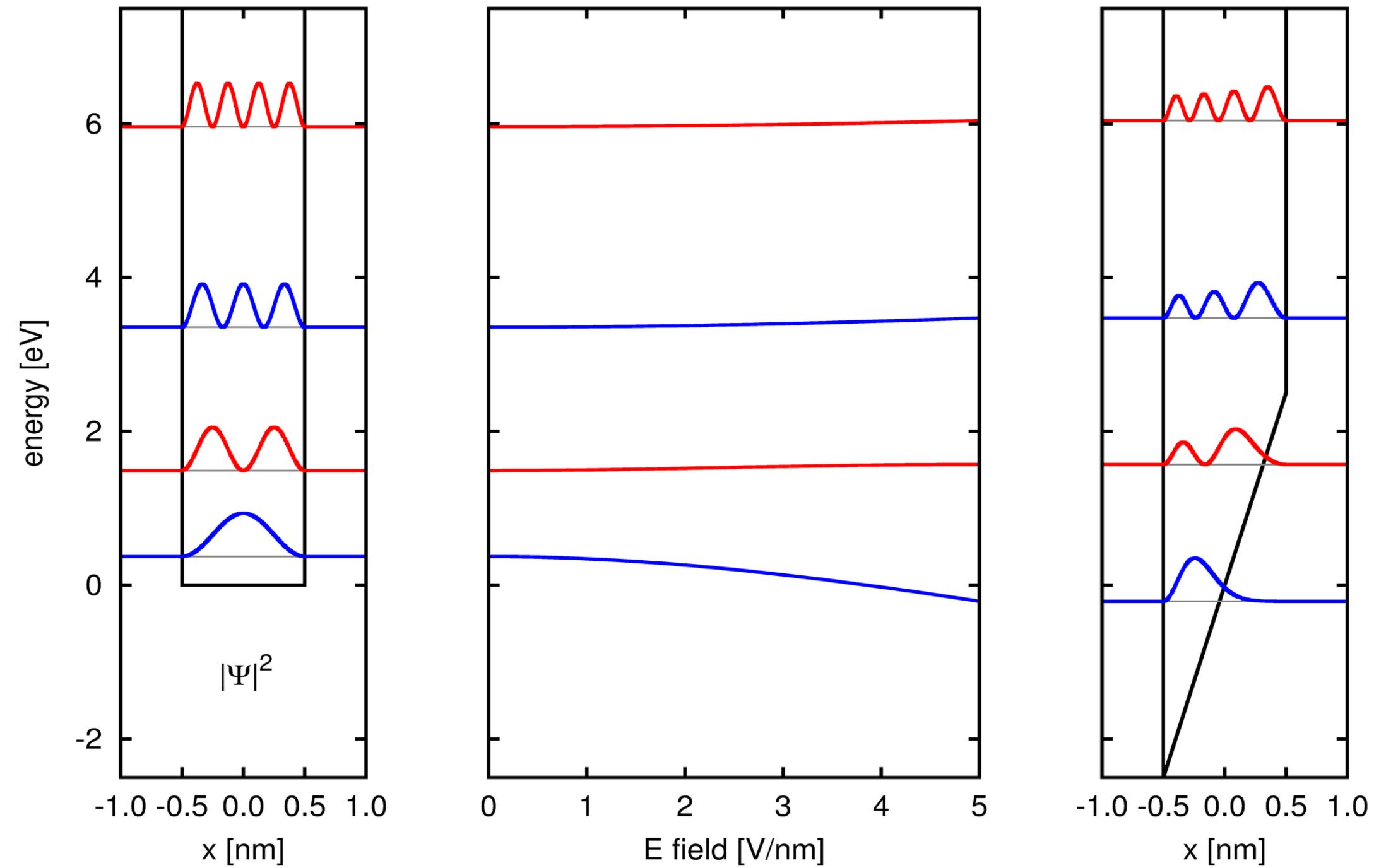
2) velikost absorpční struktury

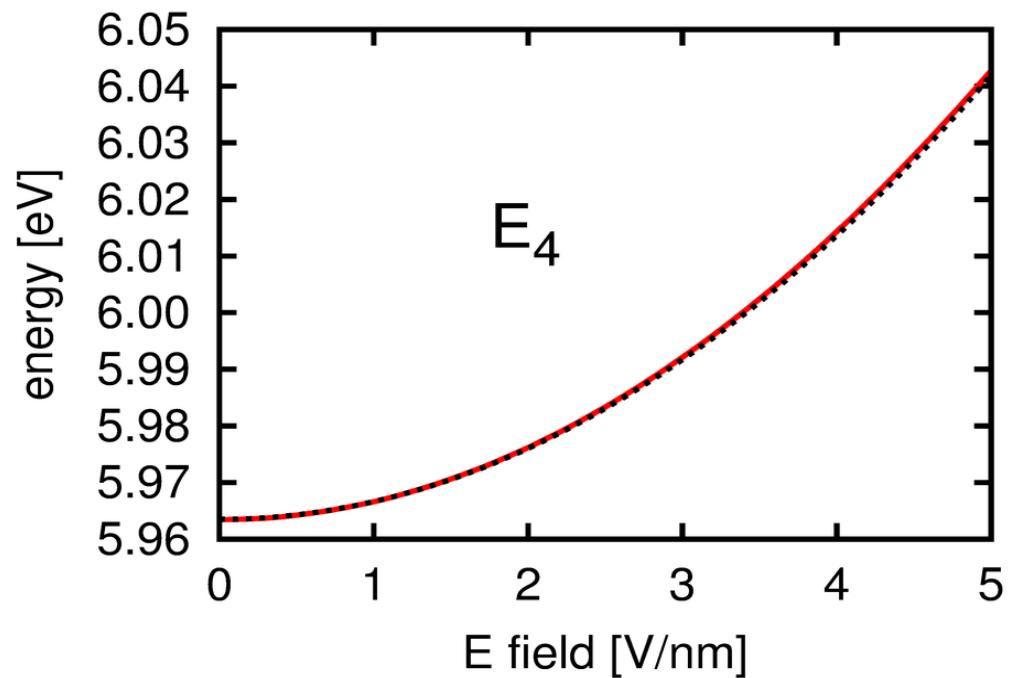
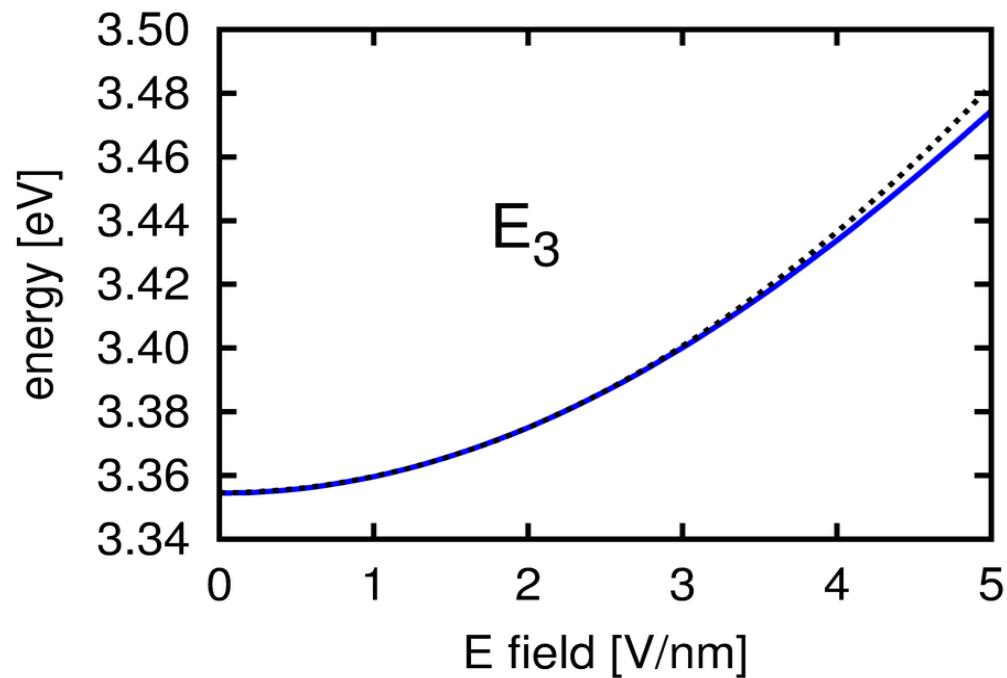
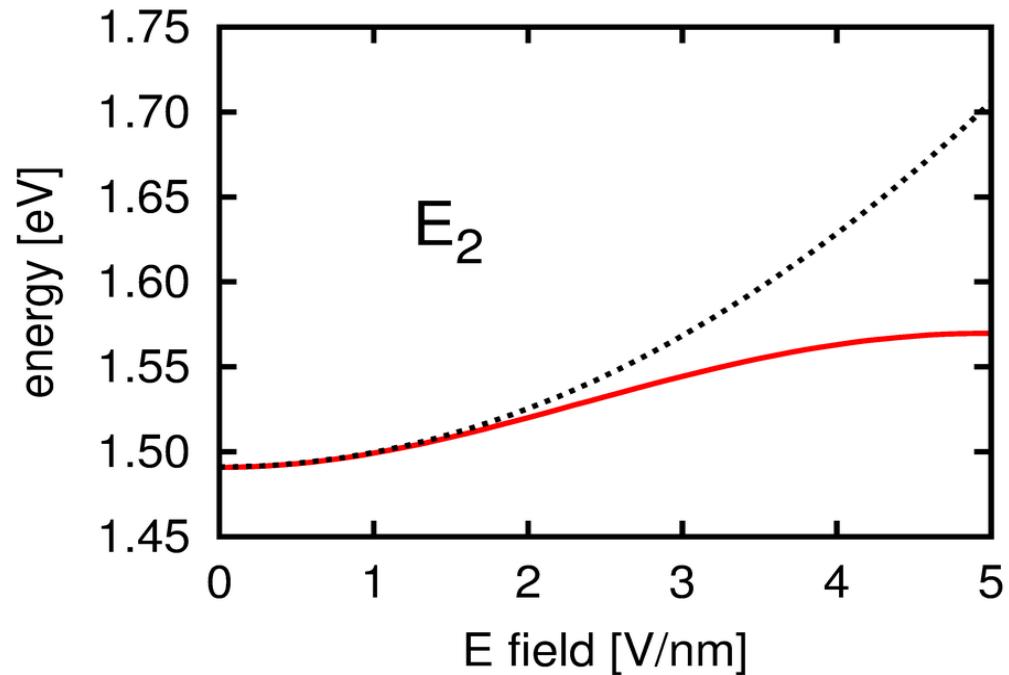
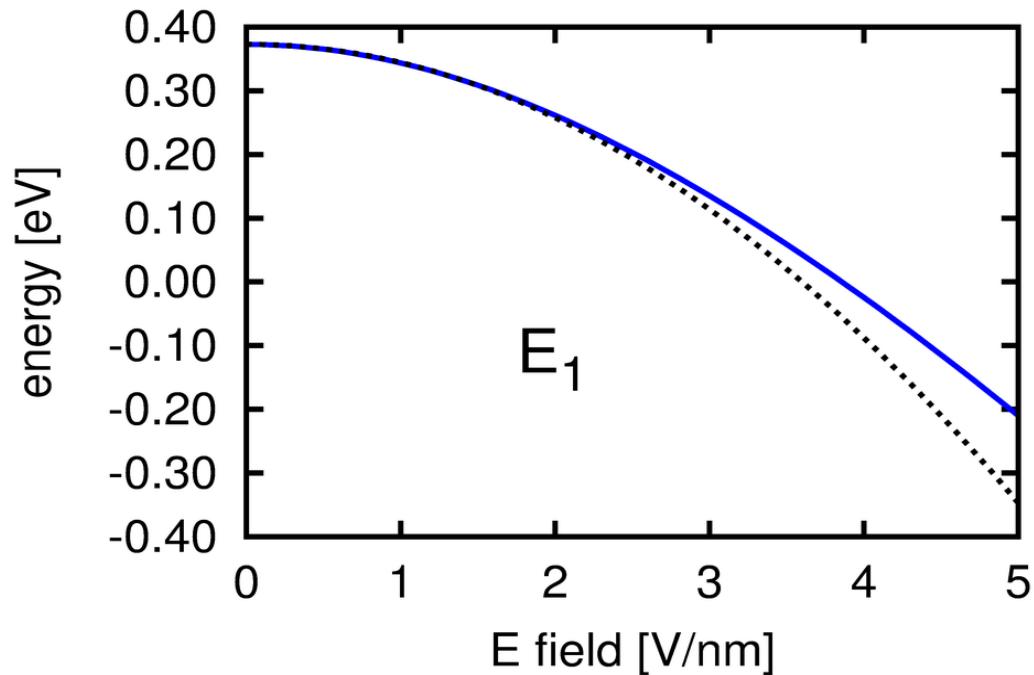
$$\sim \left| \int \psi_e^*(x) \psi_h(x) dx \right|^2 \sim \left| 1 - \frac{2(\pi^2 - 6)}{3} \lambda^2 \right|^2 \approx 1 - \frac{4(\pi^2 - 6)}{3} \lambda^2$$

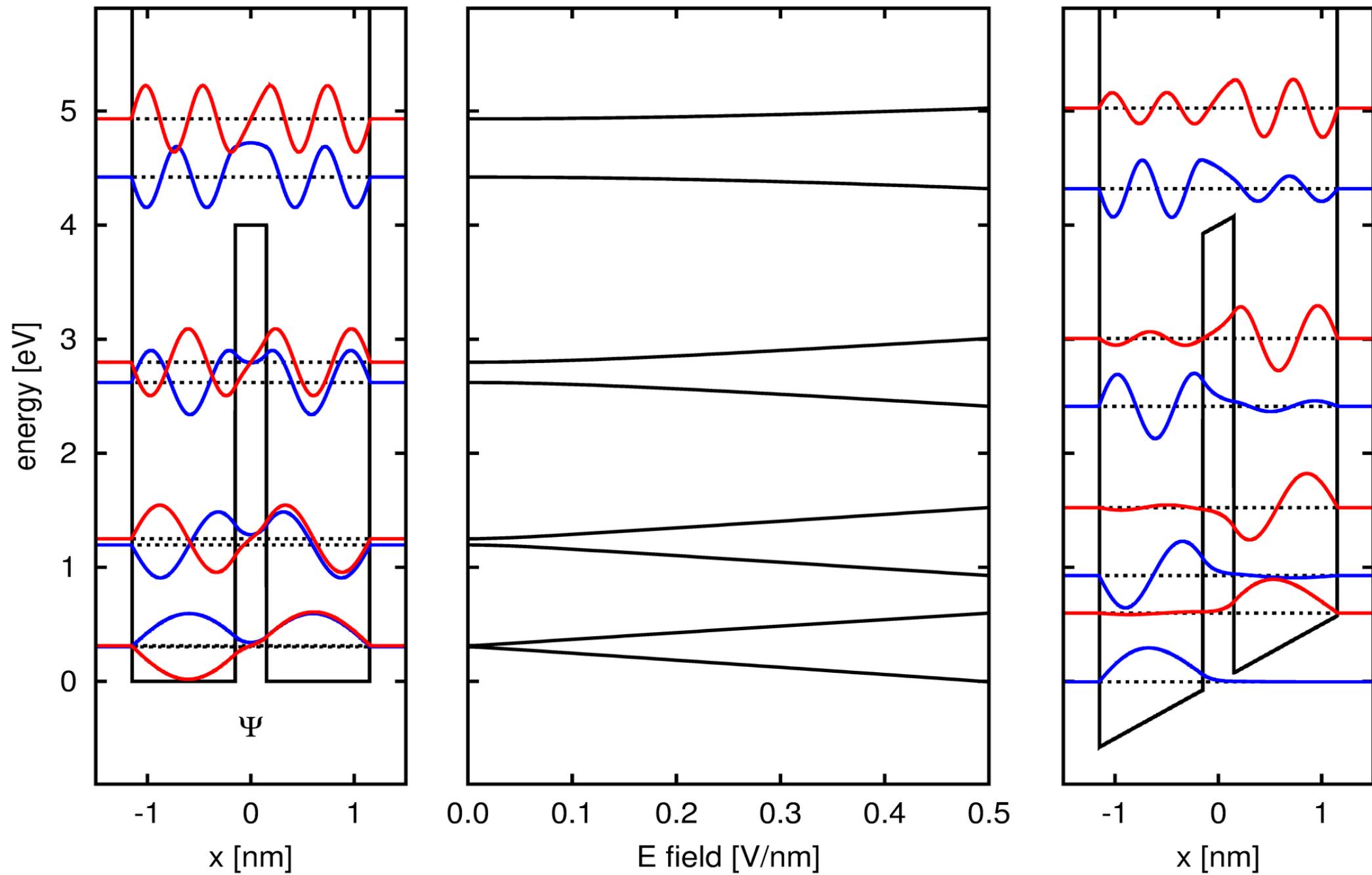
překryvoj integrál

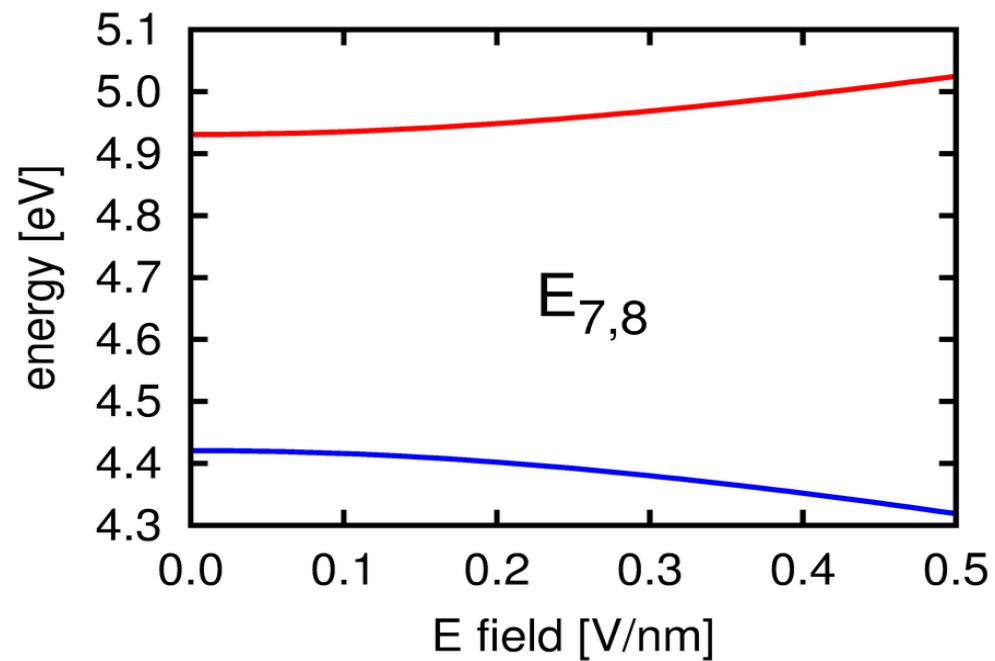
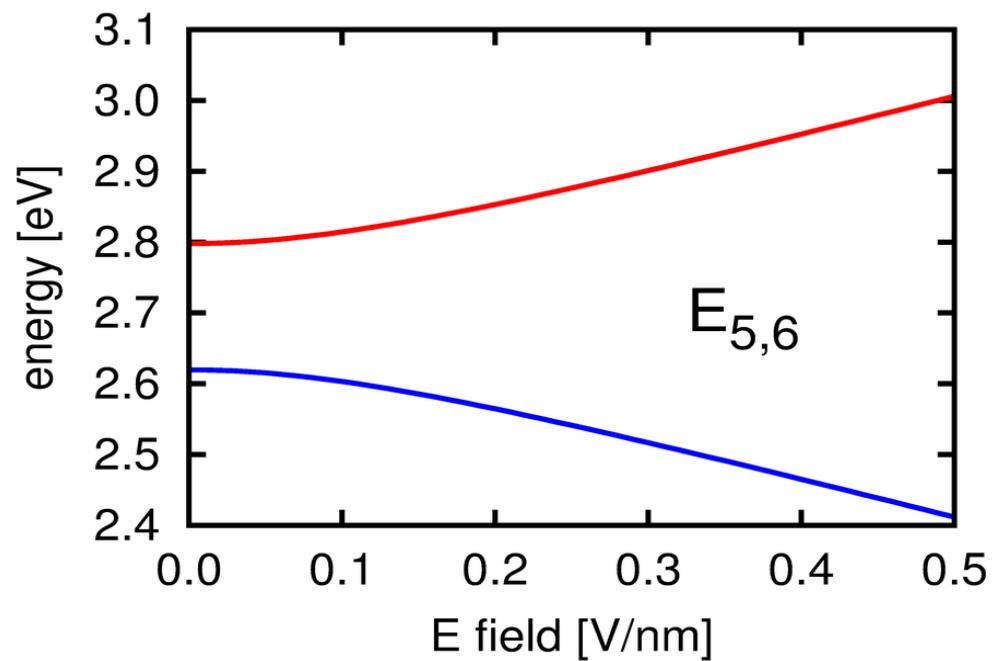
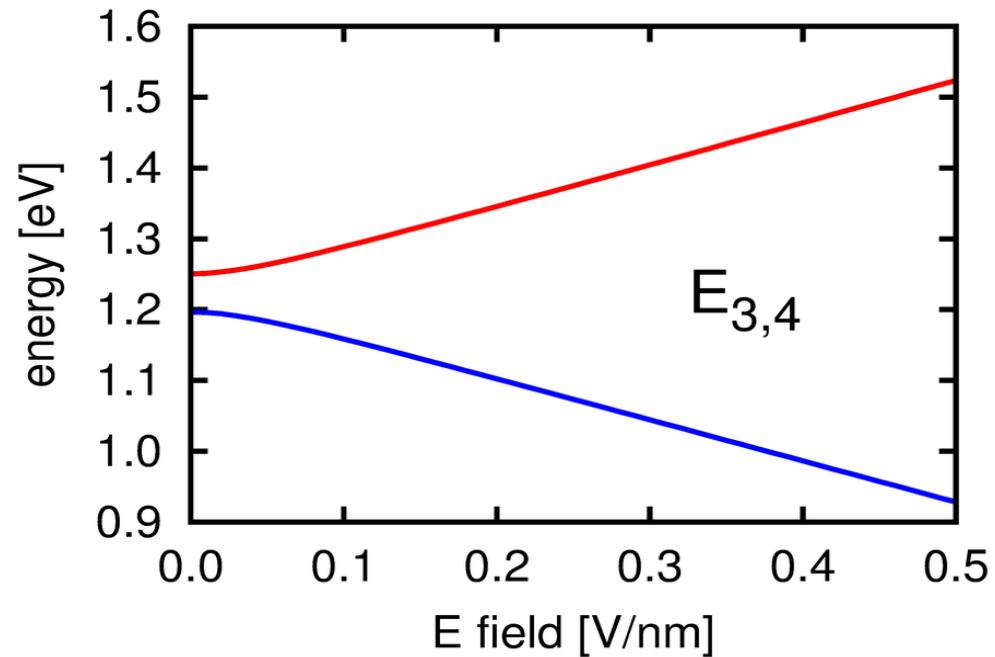
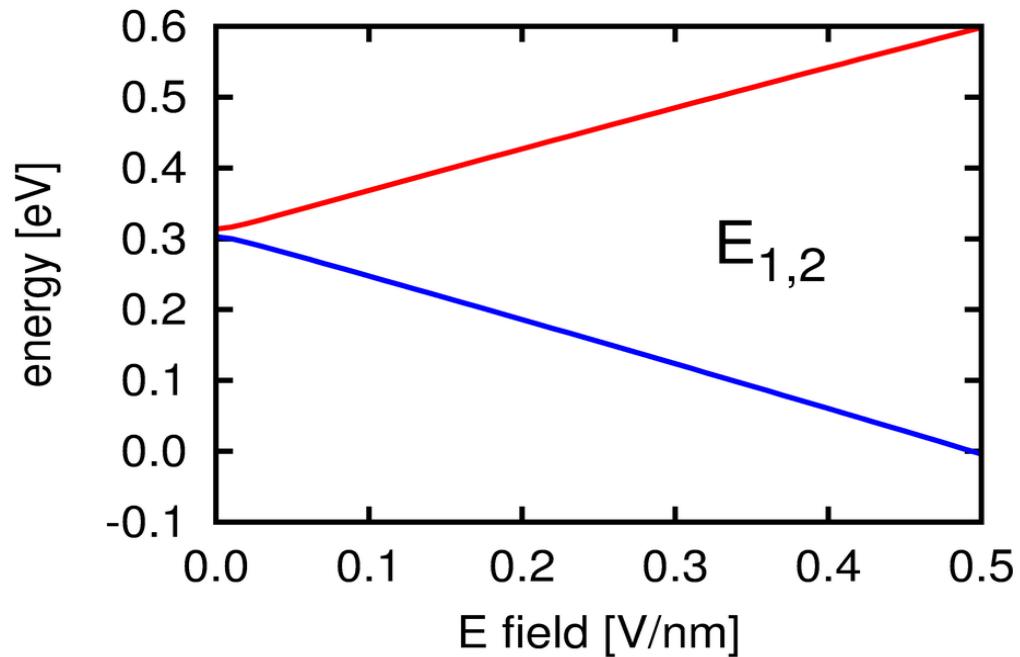
$$\lambda \sim u \rightarrow \text{úbytek absorpce} \sim u^2$$

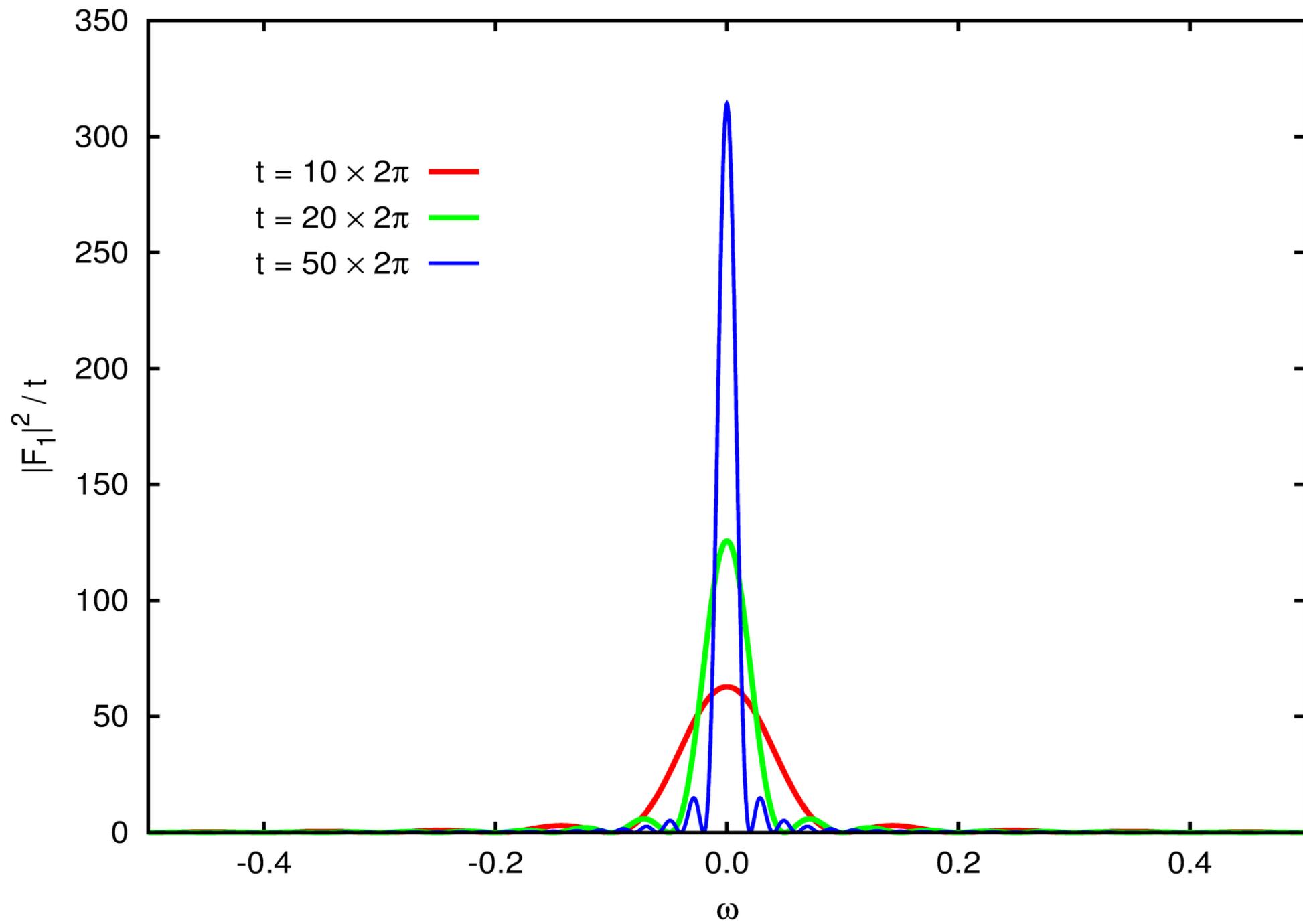


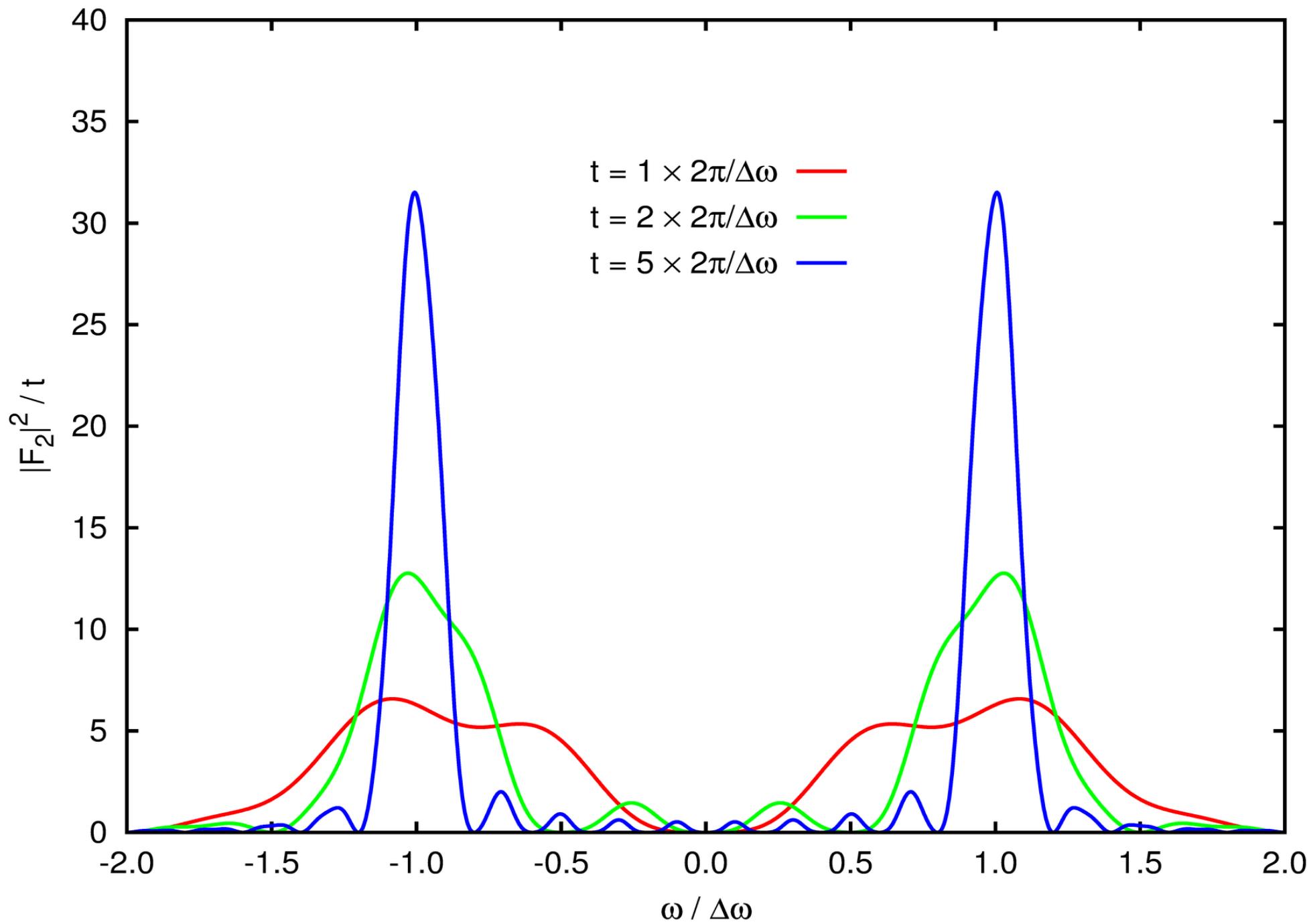


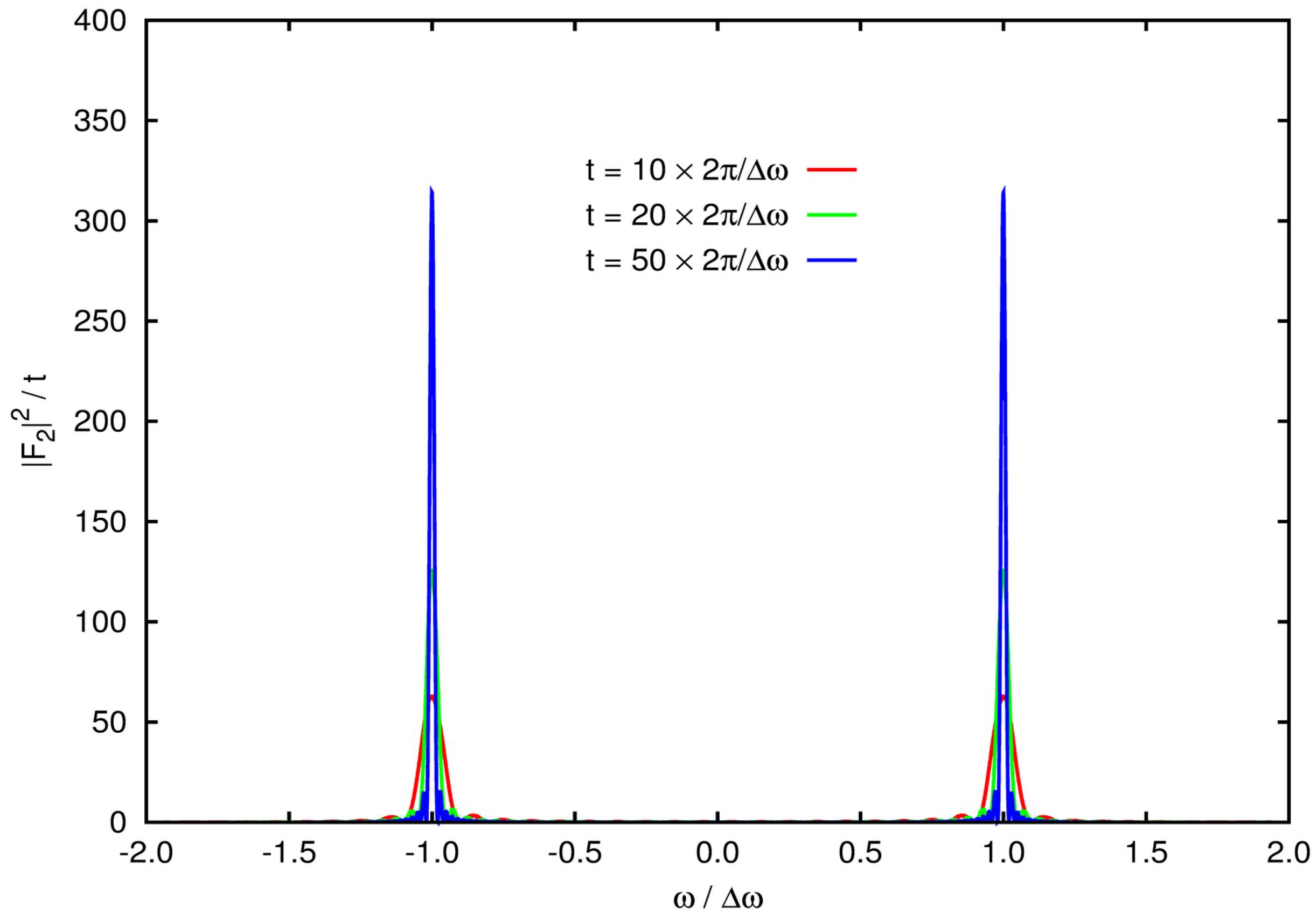


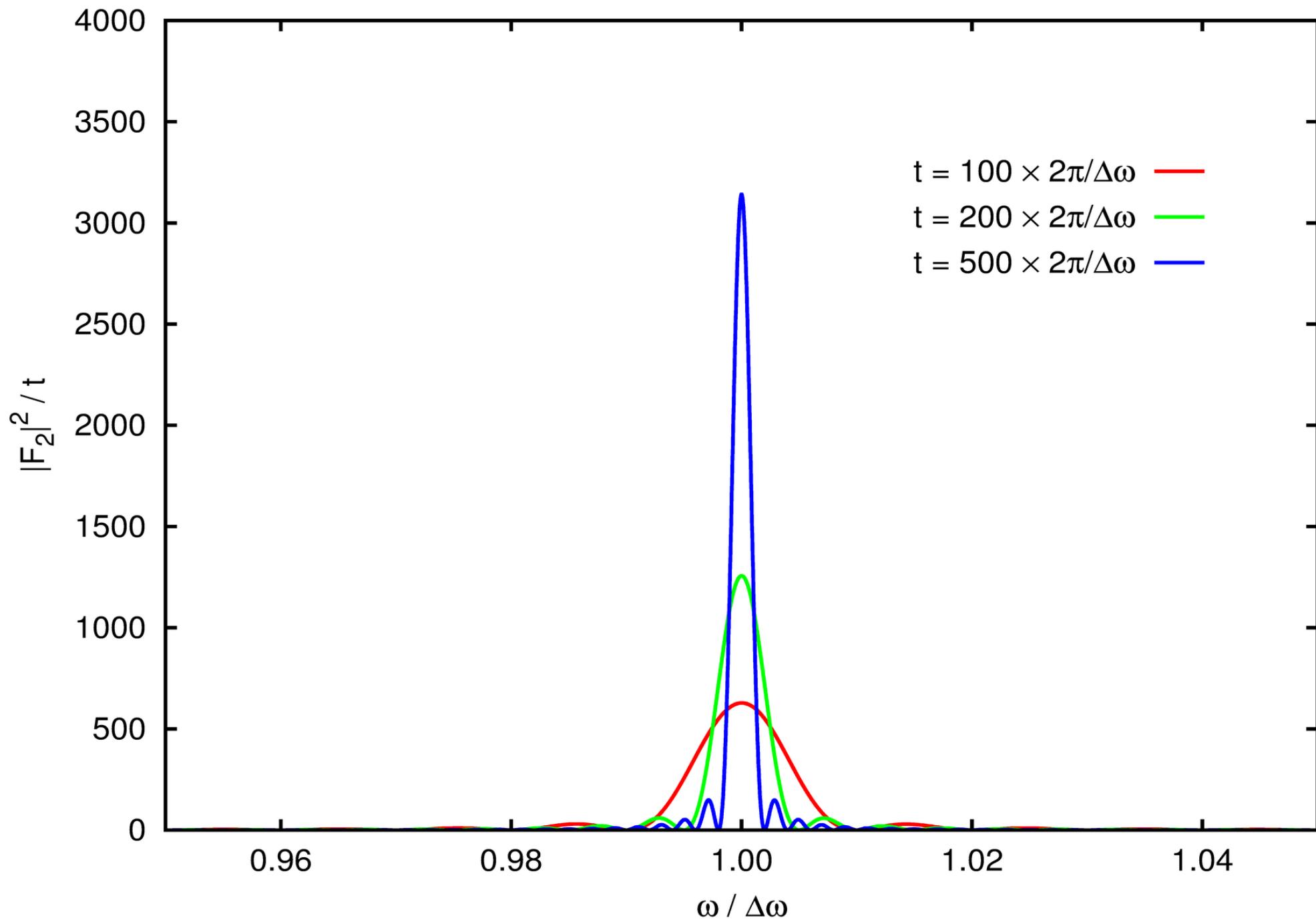


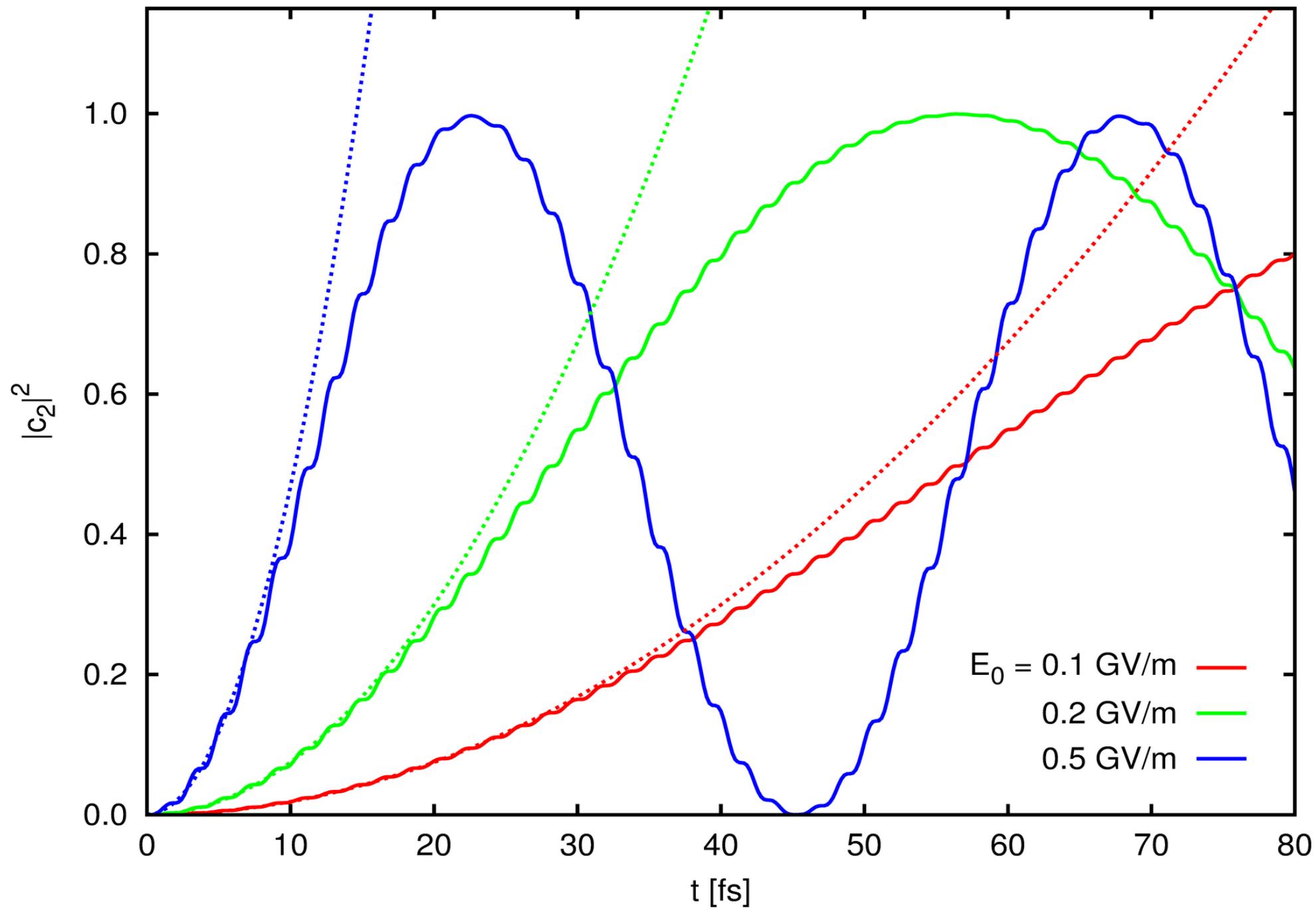


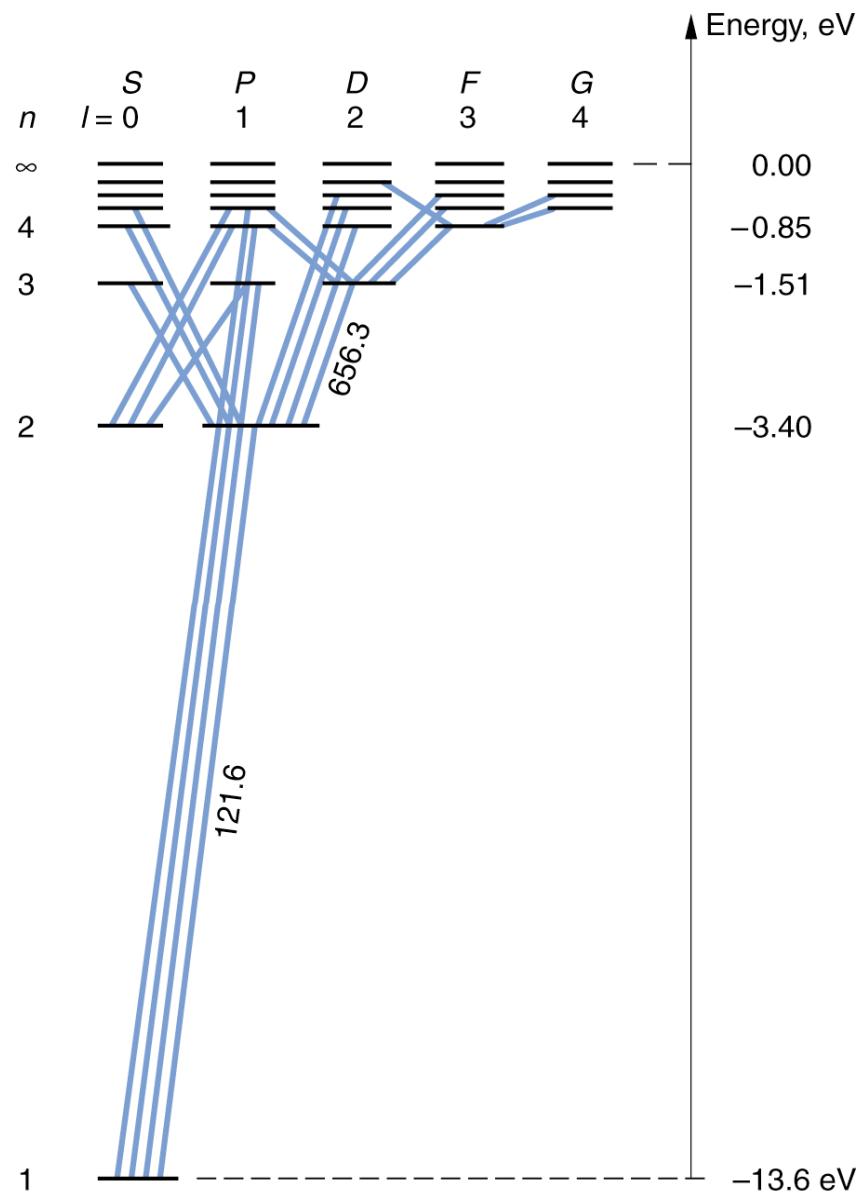
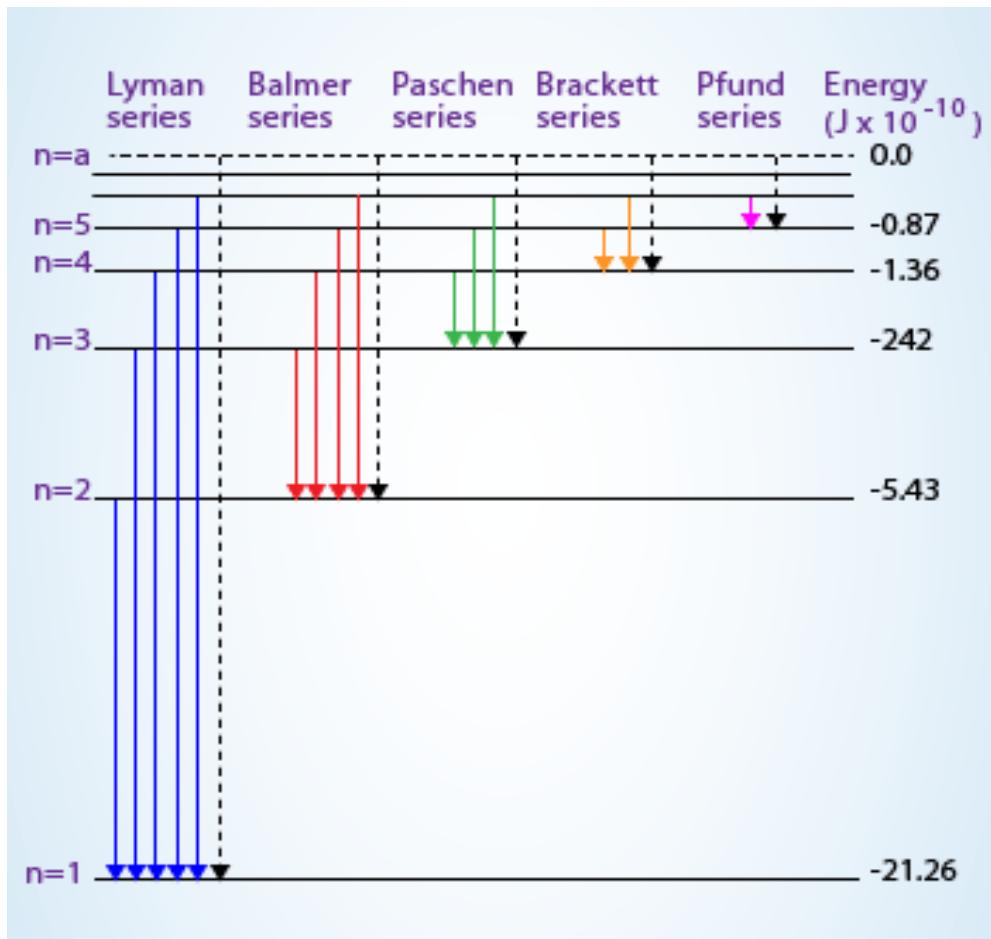


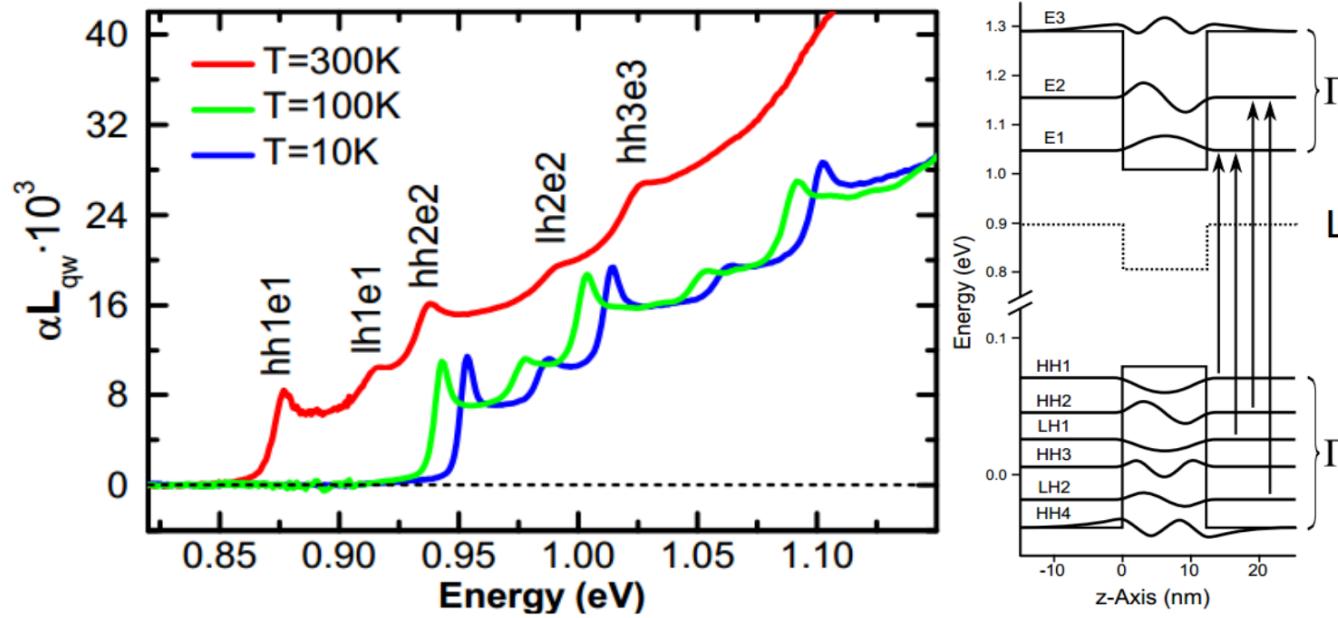












**Figure 2.** Linear absorption spectrum of the Ge multiple quantum well structure. A schematic sketch of the electronic structure according to [18] is given on the right-hand side.

