C7790 Introduction to Molecular Modelling TSM Modelling Molecular Structures

Lesson 12 Quantum Chemistry I

PS/2020 Distant Form of Teaching: Rev1

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C7790 Introduction to Molecular Modelling

Context

macroworld

states

(thermodynamic properties, G, T,...)

phenomenological thermodynamics

equilibrium (equilibrium constant) kinetics (rate constant)

free energy (Gibbs/Helmholtz)

microworld

Description levels (model chemistry):

quantum mechanics

- semiempirical methods
- ab initio methods
- post-HF methods
- DFT methods
- molecular mechanics
- coarse-grained mechanics

Simulations:

partition function

statistical thermodynamics

microstates

(mechanical properties, E)

• molecular dynamics

- Monte Carlo simulations
- docking
- ...

microstate ≠ microworld



Method overview (model chemistry]



Quantum chemistry I

Multi-electron systems

Approximations:

- Born-Oppenheimer approximation
- > One-electron approximation
- Basis functions

Chemical system

Hamiltonian of chemical system consisting of *N* nuclei of mass *M* and charge *Z* and *n* electrons is given by:



Schrödinger equation:

$$\hat{H}\psi_k(\mathbf{r},\mathbf{R}) = E_k(\mathbf{R})\psi_k(\mathbf{r},\mathbf{R})$$

Finding the solution

SE is complicated differential equation:

$$\hat{H}\psi_k(\mathbf{r},\mathbf{R}) = E_k(\mathbf{R})\psi_k(\mathbf{r},\mathbf{R})$$

(the devil is a kinetic operator)

- Instead of searching for all possible solutions, lets focus on a ground state.
- \succ The ground state is the state with the lowest energy E_0 .
- Since E_k are electronic states, it can be expected that for the most systems that the ground state will determine the essential behavior of the system.

$$E_0 < E_1 < E_2 < \dots$$

The ground state can be found be a variational method.

Alternative: a perturbation method.

Variational method

$$\hat{H}\psi_k(\mathbf{r},\mathbf{R}) = E_k(\mathbf{R})\psi_k(\mathbf{r},\mathbf{R})$$
 differential form

The **variational method** employs variational calculus. The method essence is to find the local extreme of **functional**, which is the way how to transform functions on real numbers. Such a representation is the relationship between energy and wave function expressed in integral form:

$$\int_{\Omega} \Psi^{*}(\mathbf{r},\mathbf{R})\hat{H}\Psi(\mathbf{r},\mathbf{R})d\mathbf{\tau}$$
$$E(\mathbf{R}) = \frac{\Omega}{\int_{\Omega} \Psi^{*}(\mathbf{r},\mathbf{R})\Psi(\mathbf{r},\mathbf{R})d\mathbf{\tau}}$$

integral form (fully equivalent to SE)

$$d\mathbf{\tau} = d\mathbf{r}_1 d\mathbf{r}_2 d\mathbf{r}_3 d\mathbf{r}_4 \dots d\mathbf{r}_n$$

integrates across all electrons and the entire space $\boldsymbol{\Omega}$

Variational method

$$E_{k} = E[\Psi_{k}] = \frac{\int \Psi_{k}^{*} \hat{H} \Psi_{k} d\tau}{\int \int \Psi_{k}^{*} \Psi_{k} d\tau} = \min!$$

The wave function, which provides the minimum value of the integral, is a solution of the Schrödinger equation. The global minimum of a functional is the energy of the ground state, which implies:

$$\Psi \neq \Psi_0$$
$$E > E_0$$

The inaccurate wave function always provides a higher value of energy.

Hartree-Fock method

Finding a wave function

Finding of wavefunction is practically impossible because it is a function of all electron positions. Possible simplification is **one-electron approximation**:

Hartree's method

$$\Psi(r_1, r_2, ..., r_n) = \varphi_1(r_1)\varphi_2(r_2)...\varphi_n(r_n)$$

Similar approach was employed for:

- time-independent SE
- BO approximation

Finding a wave function

Finding of wavefunction is practically impossible because it is a function of all electron positions. Possible simplification is **one-electron approximation**:

Hartree's method

$$\Psi(r_1, r_2, ..., r_n) = \varphi_1(r_1)\varphi_2(r_2)...\varphi_n(r_n)$$

Hartree's method does not consider important properties of multi-electron systems. Electrons are **indistinguishable fermions** (particles with half spin), which they must comply with **Pauli Exclusion Principle**:

no two indistinguishable fermions can be in the same quantum state

The wave function of the system must be **antisymmetric**. Antisymmetric wave functions can be obtained by all permutations between one-electron functions and spatial and spin coordinates.

Antisymmetric wave function:

$$\Psi(r_1, r_2, ..., r_n, \sigma_1, \sigma_2, ..., \sigma_n) = -\Psi(r_1, r_2, ..., r_n, \sigma_2, \sigma_1, ..., \sigma_n)$$

$$\Psi(r_1, r_2, ..., r_n, \sigma_1, \sigma_2, ..., \sigma_n) = -\Psi(r_2, r_1, ..., r_n, \sigma_1, \sigma_2, ..., \sigma_n)$$

spatial coordinates of electrons

spin coordinates of electrons (z-component of spin)

One-electron approximation

Hartree-Fock method

$$\Psi(r_1, r_2, ..., r_n, \sigma_1, \sigma_2, ..., \sigma_n) = \sum_{P} sign(P) \{\varphi_1(r_1)\chi_1(\sigma_1)\varphi_2(r_2)\chi_2(\sigma_2)...\varphi_n(r_n)\chi_n(\sigma_n)\}$$

all permutations z-component of spin (spin coordinate)

spin part of one-electron function

Alternative notation: Slater determinant

$$\Psi(r_{1}, r_{2}, ..., r_{n}, \sigma_{1}, \sigma_{2}, ..., \sigma_{n}) = \begin{vmatrix} \varphi_{1}(r_{1})\chi_{1}(\sigma_{1}) & \varphi_{2}(r_{1})\chi_{2}(\sigma_{1}) & \varphi_{n}(r_{1})\chi_{n}(\sigma_{1}) \\ \varphi_{1}(r_{2})\chi_{1}(\sigma_{2}) & \varphi_{2}(r_{2})\chi_{2}(\sigma_{2}) & \varphi_{n}(r_{2})\chi_{n}(\sigma_{2}) \\ \varphi_{1}(r_{n})\chi_{1}(\sigma_{n}) & \varphi_{2}(r_{n})\chi_{2}(\sigma_{n}) & \varphi_{n}(r_{n})\chi_{n}(\sigma_{n}) \end{vmatrix}$$

One-electron functions

- > Even with one-electron approximation, finding one-electron functions is difficult.
- Thus, one-electron functions are expressed using a linear combination of basis functions.
- This description is exact if we use complete system of basis functions (infinitely large set of orthonormal functions).
- The problem is then reduced to finding linear coefficients c, which determine extent of given basis functions to searched one-electron functions.



Basis functions

 $\varphi_i(\mathbf{r}_i) = \sum_{j=1}^m c_{ij} \chi_j(\mathbf{r}_i)$ pre-defined basis functions

From a practical (numerical) point of view, it is necessary to use **limited number of basis functions**.

Basis function choice affects **speed** of calculation and **accuracy** of achieved results.

Basis set types:

- atomic orbitals (atom centered) are usually derived from simplified SE solution for hydrogen atom
 - GTO Gaussian Type Orbital
 - STO Slater Type Orbital (more accurate but more complicated)
- plane waves solid state physics

HF method

The HF method tries to find such *c*, which minimizes energy functional.

one-electron molecular orbitals

$$\varphi_{i}(\mathbf{r}_{i}) = \sum_{j=1}^{m} c_{ij} \chi_{j}(\mathbf{r}_{i})$$
occupied orbitals
system WF

$$\Psi_{0}(r_{1}, r_{2}, ..., r_{n}, \sigma_{1}, \sigma_{2}, ..., \sigma_{n}) = \sum_{P} sign(P) \{\varphi_{1}(r_{1})\chi_{1}(\sigma_{1})\varphi_{2}(r_{2})\chi_{2}(\sigma_{2})...\varphi_{n}(r_{n})\chi_{n}(\sigma_{n})\}$$
variational method

$$E_{0} = E[\Psi_{0}] = \min!$$

RHF method

RHF - (**Restricted Hartree-Fock Method**) is applicable to closed systems (closed shell systems), where each molecular orbital contains exactly two electrons with opposite spin. Using the variational method, one-electron approximation and linear combination of basis functions, the solution can be found solving a generalized eigenproblem:

$$\mathbf{F}\mathbf{c}_i = \varepsilon_i \, \mathbf{S}\mathbf{c}_i$$

The solution is m (basis set size) of eigenvalues ε and vectors **c**. Eigenvalues ε represent the energy of one-electron functions (orbitals).



Summary

> The HF method is starting point for other QM methods

- HF is a variational method
- It uses two approximations:
 - > one-electron approximation (very bad approximation)

correlation energy (post-HF methods)

Finite number of basis functions

> correction is possible by extrapolation to a complete basis

Correlation energy



correlation energy – it is not included in the HF method because one-electron approximation

correlation energy is always negative because electron repulsion is overestimated by the HF method

Quantum chemical methods

Method classification



Quantum chemistry

time-independent Schrödinger equation

$$\hat{H}_{e}\psi_{k}(\mathbf{r},\mathbf{R})=E_{k}(\mathbf{R})\psi_{k}(\mathbf{r},\mathbf{R})$$

Formal scaling		Methods		
	HF	CI methods	MP methods	CC methods
N ⁴ -> N ² -> N ¹	HF,DFT			
N ⁵			MP2	CC2 (iterative)
N ⁶		CISD	MP3, MP4(SDQ)	CCSD (iterative)
N ⁷			MP4	CCSD(T), CC3 (iterative)
N ⁸		CISDT	MP5	CCSDT
N ⁹			MP6	
N ¹⁰		CISDTQ	MP7	CCSDTQ (iterative)

scaling, time complexity: http://en.wikipedia.org/wiki/Time_complexity

- HF Hartree–Fock method, DFT density theory functionals,
- CI configuration interaction method, MP Møller–Plesset perturbation method,
- CC coupled-clusters method, N number of basis functions

Jensen, F. Introduction to computational chemistry; 2nd ed.; John Wiley & Sons: Chichester, England; Hoboken, NJ, 2007.

QM method overview

Classification by theoretical approaches and approximations:

- empirical methods
 - extended Hückel method (EHT)

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- semi-empirical methods
 - AM1
 - PM3, PM6, PM7

• ...

- ab initio methods
 - Hartree-Fock (HF) method
 - post-HF methods
 - Møller-Plesset method (MP2, MP3, ...)
 - coupled-clusters method (CC)

• ...

- density functional theory (DFT)
 - LDA
 - GGA (BLYP, TPSS, PBE, ...)
 - hybrid (B3LYP, M06-2X, ...)

Method designation



energy and property calculations are performed on the geometry (structure), which was obtained by geometry optimization at the same level of theory

energy and property calculations geometry optimization